

Lect 20 Method of Steepest Descent

一. 最速下降法:

MSD是近似变形如

$$I(x) = \int_C h(t) e^{x p(t)} dt \quad x \rightarrow \infty$$

其中 C 为复平面上的积分线, $h(t), p(t)$ 均解析.

基本方法:

step 1: 由 $h(t), p(t)$ 的解析性将 C 变形为 C' ,
 s.t. 在 C' 上 $p(t)$ 虚部为常数.

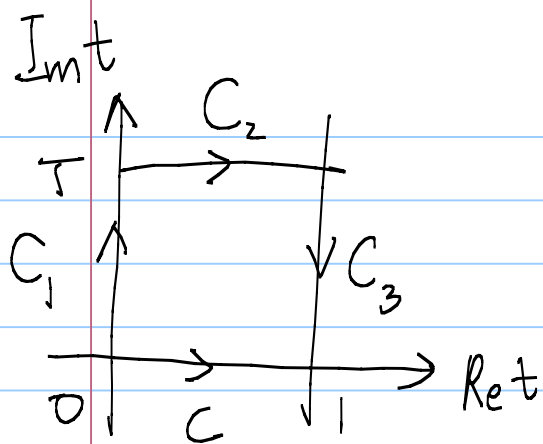
step 2: 由

$$\begin{aligned} I(x) &= \int_{C'} h(t) e^{x(\phi(t) + i\psi)} dt \\ &= e^{ix\psi} \int_{C'} h(t) e^{x\phi(t)} dt \end{aligned}$$

step 3: 由 $\phi(t)$ 为实函数, 利用 Laplace 近似.

二. 基本 MSD:

例 1: $I(x) = \int_0^1 h(t) e^{ixt} dt, \quad x \rightarrow \infty$



此时 $p(t) = it$, 取 $t = u + i\nu$ 20-02

$$p(t) = i\bar{u} - \nu$$

为使 $p(t)$ 虚部为常数即 $\nu = \text{const}$.

① $t=0$, $u=\nu=0$ 过 $t=0$ 的曲线 $t=i\nu$, $\text{Im } p(t) \equiv 0$

② $t=1$, $u=1, \nu=0$, 过 $t=1$ 的曲线 $t=1+i\nu$, $\text{Im } p(t) \equiv 1$

构造曲线 C_1, C_2, C_3 如上图, 由 Cauchy 定理:

$$I(x) = \int_C h(t) e^{ixt} dt = \int_{C_1 + C_2 + C_3} h(t) e^{ixt} dt$$

$$\text{对 } C_2: \int_{C_2} h(t) e^{ixt} dt = \int_{T\bar{i}}^{T\bar{i}+1} h(t) e^{ixt} dt$$

$$= \int_0^1 h(T\bar{i}+s) e^{ix(T\bar{i}+s)} ds$$

$$= e^{-xT} \int_0^1 h(T\bar{i}+s) e^{ixs} ds \rightarrow 0 \quad (T \rightarrow \infty)$$

$$\text{故 } I(x) = \int_{C_1} + \int_{C_3} = I_1(x) + I_3(x)$$

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$$I_1(x) = \int_0^{\infty} h(is) e^{ix(is)} i ds \quad (t = is)$$

$$= -\frac{\pi}{2x} - i \frac{h'x}{x} + \underbrace{\frac{i}{x} \int_0^{\infty} h y e^{-y} dy}_{\text{Euler 常数}}$$

$$I_3(x) = i \int_{+\infty}^0 h(1+is) e^{ix(1+is)} ds, \quad t = 1+is$$

$$= i e^{ix} \int_{+\infty}^0 h(1+is) e^{-xs} ds$$

$$\sim i e^{ix} \sum_{n=1}^{\infty} \frac{(-i)^n (n-1)!}{x^{n+1}} \quad \text{Watson 定理}$$

$$\text{故 } I(x) \sim -\frac{i h'x}{x} - \frac{i\gamma + \pi/2}{x} + i e^{ix} \sum_{n=1}^{\infty} \frac{(-i)^n (n-1)!}{x^{n+1}}$$

三. 关于最速下降路径的分析:

$$\text{对 } I(x) = \int_C e^{x f(t)} h(t) dt, \text{ 设 } f(t) = \phi(t) + i\psi(t)$$

为 $t = u + i\nu$ 的解析函数, 则 ϕ, ψ 满足 C-R 方程

$$\frac{\partial \phi}{\partial u} = \frac{\partial \psi}{\partial \nu}, \quad \frac{\partial \phi}{\partial \nu} = -\frac{\partial \psi}{\partial u}$$

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则有 $\nabla\phi, \nabla\psi = 0$

考虑过 $\vec{t}_0$ 的 $\phi(t)$ 的最速降线 $\vec{t}(s) = (u(s), v(s))$, 则

$$\begin{cases} \frac{d\vec{t}}{ds} = -\nabla\phi(\vec{t}(s)) \\ \vec{t}|_{s=0} = \vec{t}_0 \end{cases}$$

过 $\vec{t}_0$ 的常相曲线对应于 $\psi(\vec{t}(s)) = \text{const.}$

即 $\frac{d\psi(\vec{t}(s))}{ds} = 0$, 而 $\frac{d\psi(\vec{t}(s))}{ds} = \nabla\psi \cdot \frac{d\vec{t}}{ds} = -\nabla\psi \cdot \nabla\phi = 0$

即最速降线与常相曲线重合!

当 $P'(t) \neq 0$ 时 即 $P'(t) = \frac{\partial\phi}{\partial u} - i \frac{\partial\phi}{\partial v} \neq 0$ 即

$\nabla\phi \neq 0$ 时最速降线可无穷延伸, 否则如在 $\vec{t}_0$ 点停止则由 $\nabla\phi(\vec{t}_0) \neq 0 \Rightarrow \frac{d\vec{t}}{ds} \neq 0$ 矛盾! 因此最值仅在端点出现; 当 $P'(t) = 0$ 时, 称为鞍点, 此时由于右端 $= 0$ 可导致从不同初值出发会相交, 最值可

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能在曲线中间取到

例3: $P(t) = t^2$, 鞍点 $t=0$

$$\text{令 } t = u + i\nu, \quad P(t) = u^2 - \nu^2 + i2u\nu$$

零相曲线 $\psi(t) \equiv 0 \Rightarrow u=0$ 或 $\nu=0$

$u=0 \Rightarrow \phi(t) = -\nu^2$ 在 $t=0$ 最大

$\nu=0 \Rightarrow \phi(t) = u^2$ 在 $t=0$ 最小.



$$\text{例4: } P(t) = \bar{t} \cosh t = \bar{t} \cdot \frac{e^t + e^{-t}}{2}$$

$$P'(t) = 0 \Rightarrow t=0 \text{ 为鞍点}$$

$$\text{令 } t = u + i\nu, \quad P(t) = \underbrace{-\sinh u \sin \nu}_{\phi(t)} + \underbrace{\bar{t} \cosh u \cos \nu}_{\psi(t)}$$

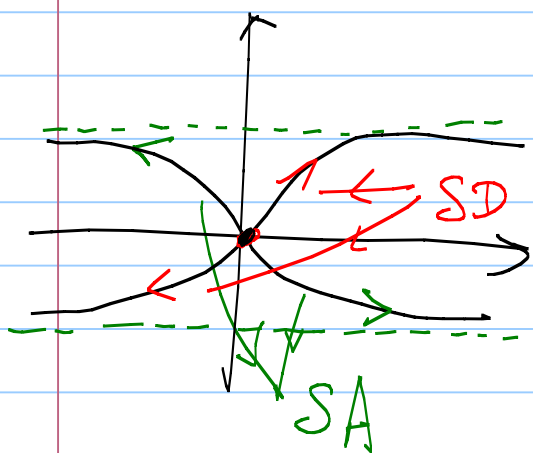
$$P(0) = \bar{t}$$

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常相曲线

$$\psi(t) = \cosh u \cos v = 1$$

$$\Rightarrow \frac{u'}{v'} = \lim_{u \rightarrow 0} \frac{\sqrt{1 - \cosh^2 u}}{\sinh u} \sim \frac{u}{u} = 1$$

事实上 $t = \pm i n \pi$ 也为鞍点

例5: $p(t) = \sinh t - t$

$$p'(t) = 0 \Rightarrow t = 0, \quad p''(0) = 0, \quad p'''(0) = 1,$$

称 $t = 0$ 为 3 阶鞍点.

令 $t = u + iv$, $p(t) = (\sinh u \cos v - u) + i(\cosh u \sin v - v)$

常相曲线 $\cosh v \sin v - v = 0$

四. 带鞍点的 SDA:

例6: $I(x) = \operatorname{Re} \frac{1}{i\pi} \int_{-i\pi/2}^{i\pi/2} e^{ix \cosh t} dt$

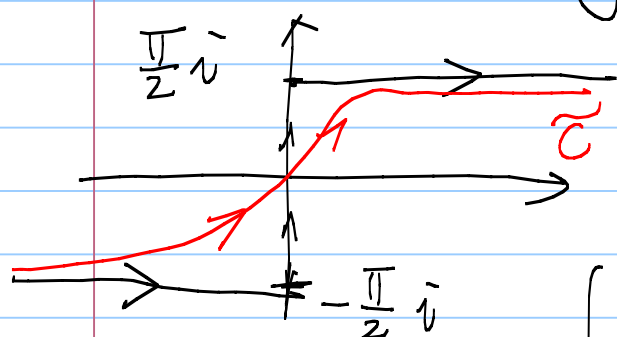
引入积分, $\frac{1}{i\pi} \int_{-\infty - i\pi/2}^{-i\pi/2} e^{ix \cosh t} dt$

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及 $\frac{1}{i\pi} \int_{i\frac{\pi}{2}}^{i\frac{\pi}{2}+\infty} e^{\bar{u}x \cosh t} dt$

可得其积分收敛且连续. 故

$$I(x) = \operatorname{Re} \frac{1}{i\pi} \int_C dt e^{\bar{u}x \cosh t}$$



由解析性及 $\pm \frac{\pi}{2}i$ 为渐近线,

$$\int_C dt e^{\bar{u}x \cosh t} = \int_{\tilde{C}} dt e^{\alpha P(t)}$$

$P(t) = \bar{u} \cosh t$, 故

$$\int_{\tilde{C}} dt e^{\bar{u}x \cosh t} = \int_{-\infty}^{+\infty} dt(s) e^{\alpha P(t(s))}$$

$t(s) = u(s) + i v(s)$, s 为 $\tilde{C}$ 的某个曲线参数.

$$\sim \int_{-\varepsilon}^{\varepsilon} \frac{dt}{ds} e^{\alpha \left\{ P(t(0)) + P'(t(0)) t'(0)s + \frac{1}{2} s^2 [P''(t(0)) (t'(0))^2 + P'(t(0)) t''(0)] \right\}} ds$$

$$\sim \int_{-\varepsilon}^{\varepsilon} (1 + i\bar{u}) e^{\alpha \left\{ \bar{u} + \frac{1}{2} s^2 [\bar{u} (1 + \bar{u})^2] \right\}} ds$$

$t = (1 + i\bar{u})s$ (局部切向)

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$$\rho(t(0)) = \bar{v}, \quad \rho''(t(0)) = \bar{v}$$

$$\sim \int_{-\varepsilon}^{\varepsilon} (1+\bar{v}) e^{i\bar{v}s} \cdot e^{-xs^2} ds \sim (1+\bar{v}) e^{i\bar{v}x} \cdot \sqrt{\frac{\pi}{x}}$$

$$\text{即 } I(x) \sim \sqrt{\frac{2}{\pi x}} \cos(x - \frac{\pi}{4}), \quad x \rightarrow \infty$$

例7: $I(x) = \int_0^1 dt \, e^{-4xt^2} \cos(5xt - xt^3), \quad x \rightarrow \infty$

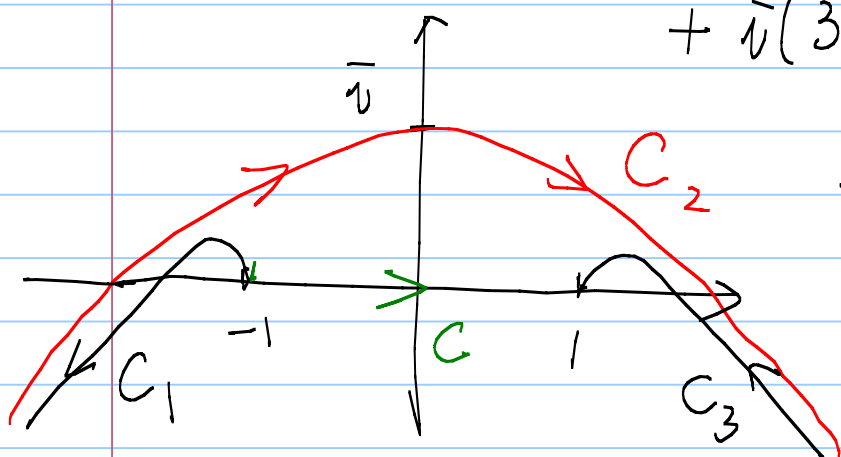
首先

$$I(x) = \frac{1}{2} \int_{-1}^1 dt \, e^{-4xt^2} \cdot e^{(5xt - xt^3) i\bar{v}}$$

$$= \frac{1}{2} e^{-2x} \int_{-1}^1 e^{x\rho(t)} dt$$

$$\rho(t) = -(t-i\bar{v})^2 - i\bar{v}(t-i\bar{v})^3$$

$$\text{令 } t = u + i\bar{v}, \quad \rho(t) = (-v^3 + 4v^2 - 5v + 3u^2v - 4u^2 + 2) + i(3u\bar{v}^2 - 8u\bar{v} + 5u - u^3)$$



过 $-1, 1$ 的常相曲线为

C_1, C_3 如左图示. 但未连接.

20-09

需要另一条曲线 C_2 将二者连接. C_2 经过鞍点 $t = \bar{v}$, 且在 $t = \bar{v}$ 处切向 $t'(s) = 1$

$$I(x) = \int_{C_1 + C_2 + C_3} \frac{1}{2} e^{-2x} e^{x p(t)} dt$$

$$\int_{C_3} e^{x p(t)} dt \sim \int_{-\varepsilon}^{\varepsilon} e^{x \left[p(t_{(0)}) + \frac{1}{2} s^2 p''(t_{(0)}) (t'(0))^2 \right]} t'(0) ds$$

$$p'(t_{(0)}) = 0, t_{(0)} = \bar{v}, t'(0) = 1, p''(t_{(0)}) = -2, p(t_{(0)}) = 0$$

$$\sim \int_{-\varepsilon}^{\varepsilon} e^{-x s^2} ds \sim \sqrt{\frac{\pi}{x}}$$

可证 $\int_{C_1}, \int_{C_3}$ 相比 $\int_{C_2}$ 为指数小.

$$\text{故 } I(x) \sim \frac{1}{2} e^{-2x} \sqrt{\frac{\pi}{x}}.$$