

The Reversible Nearest Particle System on a Finite Set *

Dayue Chen, Juxin Liu and Fuxi Zhang[†]

LMAM, School of Mathematical Sciences, Peking University, Beijing 100871, China

Abstract: In this paper we study the one-parameter family of attractive reversible nearest particle systems on $\{1, 2, \dots, N\}$. Denote by σ_N the time that the system first hits the empty set. Then, σ_N has a logarithmic increasing rate as the parameter λ is small enough, but an exponential increasing rate as λ is large enough. Especially, it has a polynomial increasing rate in the critical case, i.e. $\lambda = 1$.

Keywords: nearest particle system, first hitting time.

1 Introduction

A nearest particle system (or NPS for short) on $S \subset \mathbb{Z}$ is a spin system taking values in subsets of S , and with flip rates for any finite $A \subset S$

$$\begin{aligned} q(A, A \setminus \{x\}) &= 1 && \text{if } x \in A, \\ q(A, A \cup \{x\}) &= \beta(l_x(A), r_x(A)) && \text{if } x \in S \setminus A, \\ q(A, B) &= 0 && \text{otherwise,} \end{aligned}$$

where $l_x(A)$ and $r_x(A)$ are the distances from x to the nearest points in A to the left and right respectively. See Chapter 7 of [4] for more details.

Suppose the system is *reversible*, equivalently, by Theorem VI.1.2 of [4]

$$\beta(l, r) = \frac{\beta(l)\beta(r)}{\beta(l+r)}, \quad \beta(l, \infty) = \beta(\infty, l) = \beta(l),$$

*Supported in part by Grant G1999075106 from the Ministry of Science and Technology of China.

[†]Corresponding author

where $\sum_{l=1}^{\infty} \beta(l) < \infty$. In this paper we consider the one-parameter family $\beta_{\lambda}(l) = \lambda\psi(l)$, where $\psi(\cdot)$ is strictly positive and satisfies $\sum_{n=1}^{\infty} \psi(n) = 1$. Suppose that $\psi(n)/\psi(n+1)$ decreases to 1 as $n \rightarrow \infty$, which ensures that the process is Feller and attractive.

Denote by $\{\xi_t^N : t \geq 0\}$ the NPS on $\{1, 2, \dots, N\}$ starting from all sites occupied, and by σ_N the first time it hits the empty set. We estimate σ_N , and the results read as follows.

Theorem 1.1 *Suppose*

$$M \triangleq \sup_n \sum_{l+r=n} \frac{\psi(l)\psi(r)}{\psi(n)} < \infty. \quad (1)$$

Then, for any C_N such that $\lim_{N \rightarrow \infty} C_N = \infty$,

$$\lim_{N \rightarrow \infty} P_{\lambda}(\sigma_N \leq C_N f_{\lambda}(N)) = 1,$$

where

$$f_{\lambda}(N) = \begin{cases} \log N, & \text{if } \lambda M < 1, \\ N \log N, & \text{if } \lambda M = 1, \\ (\lambda M)^N, & \text{if } \lambda M > 1. \end{cases}$$

Theorem 1.2 *Suppose there exists n_0 such that*

$$\frac{\lambda}{\lambda'_c} > \max \left\{ \frac{2\psi(3n_0)}{\sum_{l=n_0}^{2n_0} \psi(l)\psi(3n_0-l)}, \frac{1}{\sum_{n_0}^{2n_0} \psi(l)} \right\}, \quad (2)$$

where λ'_c is the critical value for the contact process on \mathbb{Z} . Then there is a constant $\gamma > 0$ such that

$$\lim_{N \rightarrow \infty} P(\sigma_N \geq e^{\gamma N}) = 1.$$

Theorem 1.3 *Suppose $\lambda = 1$, and the initial distribution of the NPS follows the renewal measure $\text{Ren}(\beta)$ whenever the initial state is not \emptyset . Suppose C_N and C'_N are two sequences of constants such that $\lim_{N \rightarrow \infty} C_N = \infty$ and $\lim_{N \rightarrow \infty} C'_N = 0$. Then*

$$\lim_{N \rightarrow +\infty} P(C'_N N \leq \sigma_N \leq C_N N^2) = 1.$$

The initial distribution $Ren(\beta)$ of the previous theorem will be further elaborated in the beginning part of Section 4. The inequality (1) is not very restrictive. For example, let $\psi(n) = cn^{-\alpha}$, where $\alpha > 1$. A standard coupling shows that

$$\lim_{N \rightarrow \infty} P(\sigma_N > (1 - \varepsilon) \log N) = 1, \quad \forall \varepsilon > 0. \quad (3)$$

Theorem 1.1 and (3) imply that σ_N has a logarithmic increasing rate as λ is small enough. By Theorem 1.1 and Theorem 1.2, σ_N has an exponential increasing rate as λ is large enough. Theorem 1.3 tells us that σ has a polynomial increasing rate as $\lambda = 1$, the critical point of the NPS on \mathbb{Z} .

This study is inspired by a series of papers by R. Durrett *et al* [1, 2, 3], in which the contact process is concerned. Namely, let $\{\zeta_t^N : t \geq 0\}$ be the contact process on $\{1, 2, \dots, N\}$ with the parameter λ' starting from all sites occupied, and τ_N be the first time it hits the empty set. Denote by λ'_c the critical value of the contact process on \mathbb{Z} .

Theorem 1.4 (i) *If $\lambda' > \lambda'_c$, then there is a constant $\gamma_1(\lambda') > 0$ so that as $N \rightarrow \infty$, $\tau_N / (\log N) \rightarrow 1/\gamma_1(\lambda')$ in probability ([1], Theorem 1).*

(ii) *If $\lambda' > \lambda'_c$, then there is a constant $\gamma_2(\lambda') > 0$ so that as $N \rightarrow \infty$, $\log \tau_N / N \rightarrow \gamma_2(\lambda')$ in probability ([2], Theorem 2).*

(iii) *If $\lambda' = \lambda'_c$ and $a, b \in (0, \infty)$, then $P(aN \leq \tau_N \leq bN^4) \rightarrow 1$ as $N \rightarrow \infty$ ([3], Theorem 1.6).*

The contact process is a non-reversible NPS. We believe the same conclusion is also true for the reversible NPS. However, we are only able to show it for very small λ and very large λ . Moreover, the parameters in the lower estimate and the upper estimate should be amended to the same.

Theorems 1.1, 1.2 and 1.3 are proved in Sections 2, 3 and 4 in turn. In Section 3, we give a proof of (3) for the completeness.

2 Upper Estimate of σ_N

To get an upper bound of σ_N , we compare the evolution of $\{|\xi_t^N| : t \geq 0\}$ with a birth and death process on $\{0, 1, \dots, N\}$. On one hand, for any configuration

$|\xi| = i$, there are at most $i + 1$ intervals of consecutive vacant sites, which do not intersect mutually; in each interval, the rate that a new particle is born is no more than λM . Hence the rate that $|\xi_t^N|$ increases 1 is not more than $(i + 1)\lambda M$. On the other hand, when $|\xi_t^N| = i$, the rate that $|\xi_t^N|$ decreases 1 equals i , the total rate that there is a particle dying.

Let $\{X_t : t \geq 0\}$ be the birth and death process on $\{0, 1, \dots, N\}$ with death rate $a_i = i$, for any $i = 1, \dots, N$; and birth rate $b_i = (i + 1)\lambda M$, for any $i = 0, \dots, N - 1$. If initially $X_0 = x \geq |\xi^N|$, there is a coupling of $\{X_t : t \geq 0\}$ and $\{\xi_t^N : t \geq 0\}$ such that

$$P^{x, \xi^N}(X_t \geq |\xi_t^N|, \forall t \geq 0) = 1, \quad (4)$$

where P^{x, ξ^N} is the conditional distribution of the initial state (x, ξ^N) .

Proof of Theorem 1.1. Let $\tau = \inf\{t > 0 : X_t = 0\}$ be the first time that $\{X_t : t \geq 0\}$ hits 0. Let P^i be the conditional probability distribution on the initial state i , and E^i be the expectation with respect to P^i , where $i = 0, \dots, N$. By (4), σ_N is stochastically dominated by τ if $X_0 \geq |\xi_0^N|$. Therefore, for any $t \geq 0$,

$$P(\sigma_N \geq t) \leq P^N(\tau \geq t). \quad (5)$$

By the Chebyshev inequality, for any $c > 0$,

$$P^N(\tau \geq cE^N\tau) \leq \frac{E^N\tau^2}{(cE^N\tau)^2}. \quad (6)$$

It is shown in [6] that

$$E^N\tau = \sum_{i=1}^N e_i, \quad E^N\tau^2 = \sum_{i=1}^N \varepsilon_i,$$

where

$$\begin{aligned} e_i &= \frac{1}{a_i} + \sum_{k=0}^{N-2-i} \frac{b_i b_{i+1} \cdots b_{i+k}}{a_i a_{i+1} \cdots a_{i+k} a_{i+k+1}} + \frac{b_i b_{i+1} \cdots b_{N-1}}{a_i a_{i+1} \cdots a_{N-1} a_N}, \\ \varepsilon_i &= \frac{2m_i}{a_i} + \sum_{k=0}^{N-2-i} \frac{2b_i b_{i+1} \cdots b_{i+k} m_{i+k+1}}{a_i a_{i+1} \cdots a_{i+k} a_{i+k+1}} + \frac{2b_i b_{i+1} \cdots b_{N-1} m_N}{a_i a_{i+1} \cdots a_{N-1} a_N}, \end{aligned} \quad (7)$$

and $m_i = E^i\tau$ for $i = 1, \dots, N$. Notice that $m_i \leq m_N$ for any $i \leq N$. It follows that $\varepsilon_i \leq 2m_N e_i$. Therefore,

$$E^N\tau^2 = \sum_{i=1}^N \varepsilon_i \leq 2m_N \sum_{i=1}^N e_i \leq 2m_N E^N\tau = 2(E^N\tau)^2.$$

This together with (5) and (6) yields that

$$P(\sigma_N \geq c_N E^N \tau) \leq 2c_N^{-2}. \quad (8)$$

Therefore, an upper estimate of σ_N can be taken as $c_N E^N \tau$. Suppose $C_N \rightarrow \infty$ as $N \rightarrow \infty$. Let $c_N = C_N/C$, where C is the constant given in Lemma 2.1. Then the result holds by (8) and the next lemma. \square

Lemma 2.1 *There is a constant C such that for large N ,*

$$E^N \tau \leq \begin{cases} C \log N & \text{if } \lambda M < 1, \\ CN \log N & \text{if } \lambda M = 1, \\ C(\lambda M)^N & \text{if } \lambda M > 1. \end{cases}$$

Proof. By (7), for $i = 1, \dots, N$,

$$e_i \leq \begin{cases} (1 - (\lambda M)^{N-i+1}) / ((1 - \lambda M) i), & \text{if } \lambda M \neq 1; \\ (N - i + 1)/i, & \text{if } \lambda M = 1. \end{cases}$$

Hence, if $\lambda M < 1$, there is a constant C so that

$$E^N \tau = \sum_{i=1}^N e_i \leq (1 - \lambda M)^{-1} \sum_{i=1}^N i^{-1} \leq C \log N;$$

if $\lambda M > 1$, there is a constant C so that

$$E^N \tau = \sum_{i=1}^N e_i \leq (\lambda M - 1)^{-1} \sum_{i=1}^N (\lambda M)^{N-i+1} \leq C(\lambda M)^N;$$

and if $\lambda M = 1$, there is a constant C so that

$$E^N \tau = \sum_{i=1}^N e_i \leq (N + 1) \sum_{i=1}^N i^{-1} - N \leq CN \log N.$$

\square

3 Lower Estimate of σ_N

We begin this section with recalling the monotone property of spin systems. Let $\{\xi_t : t \geq 0\}$ and $\{\zeta_t : t \geq 0\}$ be two spin systems with the same state space. Suppose that whenever $\xi \leq \zeta$,

$$c_1(x, \xi) \leq c_2(x, \zeta) \quad \text{if } \xi(x) = \zeta(x) = 0,$$

and

$$c_1(x, \xi) \geq c_2(x, \zeta) \quad \text{if } \xi(x) = \zeta(x) = 1.$$

Then by Theorem III.1.5 of [4] there is a coupling such that $P^{\xi, \zeta}(\xi_t \leq \zeta_t) = 1$ for all $\xi \leq \zeta$ and all $t \geq 0$. This together with Theorem 1.4 enlightens us to compare a NPS with a contact process by the *renormalization* argument. In other words, we divide $\{1, \dots, N\}$ into some subintervals and consider the existence of particles in each interval rather than at each site.

Proof of Theorem 1.2. Given n_0 , let $L = [N/n_0]$ be the integer part of N/n_0 , and divide $\{1, 2, \dots, Ln_0\}$ into subintervals

$$I_k = \{(k-1)n_0 + 1, (k-1)n_0 + 2, \dots, kn_0\}, \quad k = 1, 2, \dots, L.$$

We compare $\{\xi_t^N : t \geq 0\}$ with a contact process $\{\zeta_t^L : t \geq 0\}$ on $\{1, \dots, L\}$, whose initial state is

$$\zeta_0^L(k) = \begin{cases} 1 & \text{if } \sum_{x \in I_k} \xi_0^N(x) \geq 1; \\ 0 & \text{otherwise.} \end{cases}$$

We claim that, by choosing carefully the infection parameter of ζ_t^L ,

$$\sum_{x \in I_k} \xi_t^N(x) \geq \zeta_t^L(k), \quad \forall t \geq 0, k = 1, \dots, L. \quad (9)$$

This can be violated only when $\sum_{x \in I_k} \xi_t^N(x) = \zeta_t^L(k)$. So ξ_t^N and ζ_t^L evolve independently until $\sum_{x \in I_k} \xi_t^N(x) = \zeta_t^L(k)$ for some k and $t > 0$. A coupling is then needed to preserve the inequality 9. There are two cases.

Case 1. $\sum_{x \in I_k} \xi_t^N(x) = \zeta_t^L(k) = 1$. Equivalently, there is only one particle in the k -th subinterval in the configuration ξ of the NPS, and the individual at site k is infected in the configuration ζ of the contact process. Because both death rates are 1, we let both particles die at the same time.

Case 2. $\sum_{x \in I_k} \xi_t^N(x) = \zeta_t^L(k) = 0$. Equivalently, there are no particles in I_k and the individual at site k of ζ is healthy. Consider birth rates of both processes.

If $k = 1$, by attractiveness of the NPS, the total birth rate in I_1 is at least $\lambda \sum_{l=n_0}^{2n_0} \psi(l)$ if there are particles in I_2 . The case $k = L$ is similar. If $1 < k < L$, the total birth rate in I_k is at least $\lambda \sum_{l=n_0}^{2n_0} \psi(l)$ if there is at least one particle in I_{k-1} and no particle in I_{k+1} , or vice versa. If there are particles in both I_{k-1} and I_{k+1} , then

the total birth rate in I_k is at least $\lambda \sum_{l=n_0}^{2n_0} \psi(l)\psi(3n_0-l)/\psi(3n_0)$. Assumption (2) implies that we can choose the infection rate λ' of the contact process ζ_t^L to satisfy the following inequality.

$$\lambda'_c < \lambda' \leq \min \left\{ \lambda \sum_{l=n_0}^{2n_0} \frac{\psi(l)\psi(3n_0-l)}{2\psi(3n_0)}, \sum_{l=n_0}^{2n_0} \lambda\psi(l) \right\}.$$

Then there is $P_{\lambda'}$, a coupling of $\{\xi_t^N : t \geq 0\}$ and $\{\zeta_t^{[N/n_0]} : t \geq 0\}$, such that for any $t \geq 0$,

$$P_{\lambda'} \left(\sum_{x \in I_k} \xi_t^N(x) \geq \zeta_t^{[N/n_0]}(k), \forall k = 1, \dots, [N/n_0] \right) = 1. \quad (10)$$

For any $t \geq 0$,

$$P(\sigma_N \geq t) = P \left(\sum_{x \in I_k} \xi_t^N(x) \neq \emptyset \right) \geq P(\zeta_t^L \neq \emptyset) \geq P(\tau_L \geq t),$$

where $\tau_L = \inf\{t : \zeta_t^L = \emptyset\}$. This together with part (ii) of Theorem 1.4 implies that

$$\liminf_{N \rightarrow \infty} P(\sigma_N \geq e^{\gamma(\lambda')L/2}) \geq \lim_{N \rightarrow \infty} P(\tau_L \geq e^{\gamma(\lambda')L/2}) = 1.$$

Let $\gamma = \gamma(\lambda')/4n_0$, then the result follows. \square

Proof of (3). To be self-contained, we give a proof of (3). Let $\{\gamma_t^N : t \geq 0\}$ be a spin system on $\{1, 2, \dots, N\}$ starting from all sites occupied, in which particles die independently with rate 1 and no new particles are born. Then there is a coupling such that $P(\gamma_t^N \leq \xi_t^N, \forall t > 0) = 1$. This implies that

$$P(\xi_t^N \neq \emptyset) \geq P(\gamma_t^N \neq \emptyset) \quad \forall t \geq 0.$$

Notice that $P(\gamma_t^N(x) = 1) \geq e^{-t}$ for any $x = 1, \dots, N$, and $\gamma_t^N(x)$ are mutually independent. So

$$P(\sigma_N \geq \alpha(N)) \geq 1 - (1 - e^{-\alpha(N)})^N, \quad \forall \alpha(N) \geq 0.$$

Choose $\alpha(N)$ such that $(1 - e^{-\alpha(N)})^N$ converges to zero as $N \rightarrow \infty$. This gives the lower estimate of σ_N . Especially, let $\alpha(N) = (1 - \varepsilon) \log N$, where $\varepsilon > 0$. Then (3) follows. \square

4 The Critical Case

Theorem 1.3 can be divided into two separate statements:

$$\lim_{N \rightarrow \infty} P(\sigma_N \leq C_N N^2) = 1; \quad (11)$$

and

$$\lim_{N \rightarrow \infty} P(C'_N N \leq \sigma_N) = 1. \quad (12)$$

The two statements will be proved by two distinct approaches. We shall compare the critical NPS $\{\xi_t^N : t \geq 0\}$ with a critical NPS on \mathbb{Z} to show (11), and compare it with a modified process to prove (12).

For any $A = \{x_0, x_1, \dots, x_k\} \subset \{1, 2, \dots, N\}$, define

$$\nu_\beta(A) = \beta(x_1 - x_0)\beta(x_2 - x_1) \cdots \beta(x_k - x_{k-1}) \sum_{l=x_0}^{\infty} \beta(l) \sum_{r=N+1-x_k}^{\infty} \beta(r).$$

Let $K_N = \sum_{A \in \mathcal{S}_N \setminus \{\emptyset\}} \nu_\beta(A)$ and $\pi(A) = \nu_\beta(A)/K_N$. Then π is the induced probability measure of the renewal measure $Ren(\beta)$ restricted on $\{1, 2, \dots, N\}$.

The critical NPS $\{\xi_t^N\}$ is a Markov process taking values in \mathcal{S}_N with jump rate

$$q(A, B) = \begin{cases} 1 & \text{if } x \in A, B = A \setminus \{x\}, \\ \beta(l)\beta(r)/\beta(l+r) & \text{if } x \notin A, B = A \cup \{x\}; \\ 0 & \text{otherwise,} \end{cases}$$

and reversible with respect to π in the sense that $\pi(A)q(A, B) = \pi(B)q(B, A)$ for $A, B \in \mathcal{S}_N$, $A \neq \emptyset, B \neq \emptyset$. Throughout this section we take π to be the initial distribution of $\{\xi_t^N\}$.

Let $\{\widetilde{\xi}_t^N : t \geq 0\}$ be a Markov process on \mathcal{S}_N , which has the same transition rates as $\{\xi_t^N : t \geq 0\}$ except that particles can be born from the empty set. Namely, denote by \tilde{q} and q respectively the transition rates of $\{\widetilde{\xi}_t^N : t \geq 0\}$ and $\{\xi_t^N : t \geq 0\}$, then

$$\tilde{q}(A, B) = \begin{cases} q(A, B) & \text{if } A \neq \emptyset, \\ q & \text{if } A = \emptyset \text{ and } |B| = 1, \\ 0 & \text{otherwise,} \end{cases}$$

where $q > 0$ is a constant. Let

$$\nu_\beta(\{\emptyset\}) = q^{-1}, \quad \nu_\beta|_{\emptyset^c} = Ren(\beta)|_{\emptyset^c}.$$

Then $\tilde{\pi} = \nu_\beta / (K_N + q^{-1})$ is a reversible distribution of $\{\tilde{\xi}_t^N : t \geq 0\}$.

Proof of Equation (12). Let

$$\tau = \inf\{t \geq 0, \tilde{\xi}_t^N = \emptyset\},$$

\tilde{P} be the distribution of $\{\tilde{\xi}_t^N : t \geq 0\}$ with initial distribution π , and \tilde{E} be the expectation with respect to \tilde{P} . Notice that $\{\tilde{\xi}_t^N : t \geq 0\}$ is stationary under \tilde{P} . For any $t > 0$,

$$2t\pi(\{\emptyset\}) = \tilde{E} \int_0^{2t} 1_{\{\tilde{\xi}_s^N = \emptyset\}} ds.$$

By the Strong Markovian Property, the right side above equals

$$\begin{aligned} & \tilde{E} \tilde{E} \left(\int_0^{2t} 1_{\{\tilde{\xi}_s^N = \emptyset\}} ds \middle| \mathcal{F}_\tau \right) \geq \tilde{E} \tilde{E} \left(1_{\{\tau < t\}} \int_0^{2t} 1_{\{\tilde{\xi}_s^N = \emptyset\}} ds \middle| \mathcal{F}_\tau \right) \\ & \geq \tilde{E} \tilde{E} \left(1_{\{\tau < t\}} \int_\tau^{\tau+t} 1_{\{\tilde{\xi}_s^N = \emptyset\}} ds \middle| \mathcal{F}_\tau \right) = \tilde{P}(\tau < t) \tilde{E} \left(\int_0^t 1_{\{\tilde{\xi}_s^N = \emptyset\}} ds \middle| \tilde{\xi}_0^N = \emptyset \right) \end{aligned}$$

Denote by σ the first time $\{\tilde{\xi}_t^N : t \geq 0\}$ jumps. Then

$$\tilde{E} \left(\int_0^t 1_{\{\tilde{\xi}_s^N = \emptyset\}} ds \middle| \tilde{\xi}_0^N = \emptyset \right) \geq \tilde{E} \left(\sigma 1_{\{\sigma \leq t\}} \middle| \tilde{\xi}_0^N = \emptyset \right) = \int_0^t s \tilde{q}_\emptyset e^{-\tilde{q}_\emptyset s} ds,$$

where $\tilde{q}_\emptyset = \sum_\xi \tilde{q}(\emptyset, \xi) = Nq$. Hence

$$\tilde{P}(\tau < t) \leq 2t\pi(\{\emptyset\}) \left(\int_0^t s \tilde{q}_\emptyset e^{-\tilde{q}_\emptyset s} ds \right)^{-1} = \frac{2tq^{-1}}{K_N + q^{-1}} \left(\int_0^t Nqse^{-Nqs} ds \right)^{-1}. \quad (13)$$

Notice that

$$\begin{aligned} \tilde{P}(\tau < t) & \geq \tilde{P}(\tau < t, \tilde{\xi}_0^N \neq \emptyset) = \tilde{P}(\tilde{\xi}_0^N \neq \emptyset) \tilde{P}(\tau < t | \tilde{\xi}_0^N \neq \emptyset) \\ & = P(\sigma_N < t) K_N / (K_N + q^{-1}). \end{aligned}$$

This together with (13) yields that

$$P(\sigma_N < t) \leq 2tK_N^{-1} \left(q \int_0^t Nqse^{-Nqs} ds \right)^{-1} = 2NtK_N^{-1} \left(\int_0^{Nqt} se^{-s} ds \right)^{-1},$$

Notice that the right side does not depend on q , which is arbitrary. Let $q \rightarrow \infty$, then it follows that

$$P(\sigma_N < t) \leq 2NtK_N^{-1}, \quad \forall t > 0.$$

By (14), the lower estimate of σ_N is such t_N that t_N/N converges to zero, which implies the result. \square

Lemma 4.1 *Let*

$$S_N(x, y) = \{\xi^N : \xi^N(x) = \xi^N(y) = 1, \xi^N(z) = 0, \forall z < x, \text{ or } z > y\}.$$

Then $\nu_\beta(S_N(x, y)) \leq 1$ and there is a constant $C > 0$ such that $\nu_\beta(S_N(x, y)) \geq C$ whenever $y - x$ is large enough.

On the other hand, when N is large enough,

$$K_N = \sum_{\xi \in \mathcal{S}_N \setminus \{\emptyset\}} \nu_\beta(\xi) \geq \sum_{x=0}^{\lfloor N/3 \rfloor} \sum_{y=\lfloor 2N/3 \rfloor}^N \nu_\beta(S_N(x, y)) \geq CN^2. \quad (14)$$

where $\mathcal{S}_N = \{0, 1\}^{\{1, 2, \dots, N\}}$.

Proof of Lemma 4.1. Let X_n be the time until the first renewal $\geq n$. Then $\{X_n : n \geq 0\}$ is a Markov chain with transition probability $p(0, n) = \beta(n + 1)$, $p(n + 1, n) = 1$ for all $n \geq 0$. Since $\mu := \sum_{n=1}^{\infty} n\beta(n) < \infty$, $\{X_n : n \geq 0\}$ has an invariant distribution π , and $\pi(0) = 1/\mu$. Thus $P(X_n = 0)$ converges to $1/\mu$ as $n \rightarrow +\infty$. Notice that $X_n = 0$ if and only if n is a renewal time. Hence, there exists $n_0 > 0$ such that $P(X_n = 0) > (2\mu)^{-1}$ for any $n \geq n_0$. Then

$$\nu_\beta(S_N(x, y)) = P(X_{y-x} = 0) > (2\mu)^{-1}, \quad \text{if } y - x > n_0.$$

It is not difficult to check $\nu_\beta(S_N(x, y)) = P(X_{y-x} = 0) \leq 1$. □

Let $\{\eta_t : t \geq 0\}$ be a reversible nearest particle system on \mathbb{Z} , and r_t the rightmost particle in $\{\eta_t : t \geq 0\}$, i.e.

$$r_t := \sup\{x : \eta_t(x) = 1\}.$$

The properties of r_t of the critical NPS are studied in [5].

Lemma 4.2 ([5], Theorem 1) *Let $\{\eta_t : t \geq 0\}$ be the critical reversible nearest particle system on \mathbb{Z} . Suppose the initial configurations have a particle at the origin and no particle on the left of the origin, and follows the renewal measure $\text{Ren}(\beta)$ with density $\beta(\cdot)$. Then, as $a \rightarrow \infty$, r_{a^2t}/a converges in distribution to a Brownian motion with diffusion constant $D > 0$ in the Skorohod space.*

Proof of Equation (11). Since the transition mechanism of $\{\eta_t : t \geq 0\}$ is translation invariant, we can regard the NPS on $[0, N]$ as the NPS on $[n, N + n]$ for any n . So we do not distinguish them in symbols.

To use Lemma 4.2, we partition the configurations of $\{0, 1\}^N$ by the position of the rightmost particle. Namely, let A_x be the set of configurations whose rightmost particle is at x , i.e.

$$A_x = \{\xi \in \{0, 1\}^{\{0, \dots, N\}} : \xi(x) = 1, \xi(y) = 0, \forall y > x.\}$$

Denote by P be the distribution of $\{\xi_t^N : t \geq 0\}$ with initial distribution in Theorem 1.3, and by $P_{N,x}$ the distribution of the NPS on $[-N, \dots, 0]$ whose initial configurations have a particle at x , no particle to the right of x , and follows the renewal measure $Ren(\beta)$. Then

$$P = \sum_{x=0}^N P(A_x)P_{N,x}. \quad (15)$$

Denote by \mathbf{P} the distribution of the NPS on \mathbb{Z} with the initial distribution in Lemma 4.2. Now regard $\mathbf{P}_{N,x}$ the distribution of the NPS on $[-x, N-x]$. Thanks to the attractive property, there is a coupling of \mathbf{P} and $\tilde{P}_{N,x}$ such that for all $t > 0$,

$$\xi_t^N(i) \leq \eta_t(i), \quad -x \leq i \leq N-x \quad (16)$$

almost surely if $\eta_0|_{[-x,0]} = \xi_0^N$. By (16), $\xi_t^N \equiv \emptyset$ once $r_t < -x$, hence $\sigma_N \leq \inf\{t : r_t < -x\}$ almost surely.

Suppose $\lim_{N \rightarrow \infty} C_N = \infty$. For any $C > 0$,

$$\begin{aligned} P_{N,x}(\sigma_N \leq C_N N^2) &\geq P_{N,x}(\sigma_N \leq C_N x^2) \geq \mathbf{P}(\exists t \leq C_N x^2 \text{ s.t. } r_t < -x) \\ &\geq \mathbf{P}(\exists t \leq C_N \text{ s.t. } r_{x^2 t} < -x) \geq \mathbf{P}(\exists t \leq C \text{ s.t. } r_{x^2 t}/x < -1), \end{aligned}$$

whenever N is large. This together with Lemma 4.2 implies that

$$\liminf_{N,x \rightarrow +\infty} P_{N,x}(\sigma_N \leq C_N N^2) \geq \mathbf{P}(\exists t \leq C \text{ s.t. } B_t < -1), \quad \forall C > 0,$$

where $\{B_t : t \geq 0\}$ is a Brownian motion with diffusion constant D . Let $C \rightarrow +\infty$, the right side of the above equation converges to 1. Hence

$$\lim_{N,x \rightarrow +\infty} P_{N,x}(\sigma_N \leq C_N N^2) = 1.$$

Consequently, for any $\varepsilon > 0$, there exists $N_0 > 0$ such that

$$P_{N,x}(\sigma_N \leq C_N N^2) > 1 - \varepsilon.$$

for any $N \geq x \geq N_0$. This together with (15) implies that

$$P(\sigma_N \leq C_N N^2) = \sum_{x=0}^N P(A_x) P_{N,x}(\sigma_N \leq C_N N^2) \geq (1 - \varepsilon) \sum_{x=N_0}^N P(A_x). \quad (17)$$

By Lemma 4.1, on one hand,

$$\sum_{x=0}^{N_0-1} \nu_\beta(A_x) \leq \sum_{x=0}^{N_0-1} \sum_{y=0}^x \nu_\beta(S_N(y, x)) \leq N_0^2.$$

Therefore, as $N \rightarrow \infty$,

$$\sum_{x=N_0}^N P(A_x) \geq 1 - N_0^2/(C_N N^2) \rightarrow 1.$$

This together with (17) implies that

$$\liminf_{N \rightarrow \infty} P(\sigma_N \leq C_N N^2) \geq (1 - \varepsilon).$$

Let $\varepsilon \rightarrow 0$ and the result follows. □

References

- [1] Durrett, R. and Liu, X. F. (1988). The contact process on a finite set. *Ann. Probab.* **16** 1158–1173.
- [2] Durrett, R. and Schonmann, R. H. (1988). The contact process on a finite set II. *Ann. Probab.* **16** 1570–1583.
- [3] Durrett, R., Schonmann, R. H. and Tanaka, N. I. (1989). The contact process on a finite set III: The critical case. *Ann. Probab.* **17** 1303–1321.
- [4] Liggett, T. M. (1985). *Interacting particle systems*. Springer-Verlag, New York.
- [5] Schinazi, R. (1992). Brownian fluctuations of the edge for critical reversible nearest particle systems. *Ann. Probab.* **20** 194–205.
- [6] Wang Z. K. (1980). *Birth and Death process and Markov Chains* (in Chinese). Science Publishing House.

E-mail: dayue@math.pku.edu.cn

jxliu@pku.edu.cn

zhangfxi@math.pku.edu.cn

Last revision: April 17,2004