# Lecture 21. Application in Chemical Reaction Kinetics

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#### Mathematical Setup

Stochastic Simulation Algorithm (SSA)

#### Tau-leaping Algorithm

Avoiding Negative Populations Stiff System Convergence Analysis Stationary Distribution

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#### Traditional chemical reaction dynamics — ODE

Decaying-dimerizing reaction

$$\begin{array}{cccc} S_1 & \longrightarrow & \emptyset, \\ S_1 + S_1 & \longleftrightarrow & S_2, \\ S_2 & \longrightarrow & S_3, \end{array}$$

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 Traditional model — ODEs for the concentration (Law of Mass Action)

$$\frac{dx_1}{dt} = -k_1 x_1 - 2k_2 x_1^2 + 2k_3 x_2$$
$$\frac{dx_2}{dt} = k_2 x_1^2 - k_3 x_2 - k_4 x_2$$
$$\frac{dx_3}{dt} = k_4 x_2$$

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 $k_1, k_2, k_3, k_4$  are reaction rates.

### Drawbacks of ODE description

 Deterministic model describes an average behavior and is valid for large population

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Examples of stochasticity

 A. Arkin et al., Genetics 149 (1998), 1633 — Stochastic variations can produce probabilistic pathway selection.
 M. Elowitz et al., Science 297 (2002), 391 — Gene expression is affected by both extrinsic and intrinsic noise.

# Chemical kinetic system (CKS)

Taking into account the stochasticity in biological chemical reactions, this opens a new way for modeling and simulation!

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Chemical reaction kinetics - stochastic version

► Well-stirred system of N molecular species {S<sub>1</sub>, S<sub>2</sub>,..., S<sub>N</sub>} interacting through M chemical reaction channels {R<sub>1</sub>, R<sub>2</sub>,..., R<sub>M</sub>}.

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▶ Here  $a_j(x)dt$  gives the probability that the system will experience an  $R_j$  reaction in the next infinitesimal time dtwhen the current state  $X_t = x$ .  $\nu_j^i$  is the change in the number of  $S_i$  molecules caused by one  $R_j$  reaction.

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• We will define 
$$a_0(\boldsymbol{x}) = \sum_{j=1}^M a_j(\boldsymbol{x}).$$

#### An example

Decaying-dimerizing reaction:

$$S_1 \xrightarrow{R_1} \phi$$

$$S_1 + S_1 \xrightarrow{R_2} S_2$$

$$S_1 + S_1 \xleftarrow{R_3} S_2$$

$$S_2 \xrightarrow{R_4} S_3$$

The propensity functions are given by

$$a_1 = x_1, \ a_2 = 5x_1(x_1 - 1), \ a_3 = 1000x_2, \ a_4 = 0.1x_2$$

and state change vector

 $\nu_1 = (-1, 0, 0), \ \nu_2 = (-2, 1, 0), \ \nu_3 = (2, -1, 0), \ \nu_4 = (-1, 0, 1).$ Initial state X(0) = (400, 798, 0).

#### Chemical Master Equation

The Chapman-Kolmogorov equation in the time interval [t, t + dt)

$$P(x, t + dt | x_0, t_0) = \sum_{j=1}^{M} P(x - \nu_j, t | x_0, t_0) a_j(x - \nu_j) dt +$$
  
 $\left(1 - \sum_{j=1}^{M} a_j(x) dt\right) P(x, t | x_0, t_0),$ 

where dt is an infinitesimal time, and we have already omitted the higher order terms in o(dt).

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where dt is an infinitesimal time, and we have already omitted the higher order terms in o(dt).

With some algebra, we get the well-known chemical master equation.

$$\partial_t P(\boldsymbol{x}, t | \boldsymbol{x}_0, t_0) = \sum_{j=1}^M a_j(\boldsymbol{x} - \boldsymbol{\nu}_j) P(\boldsymbol{x} - \boldsymbol{\nu}_j, t | \boldsymbol{x}_0, t_0) - a_0(\boldsymbol{x}) P(\boldsymbol{x}, t | \boldsymbol{x}_0, t_0)$$

where  $a_0(x):=\sum_{j=1}^M a_j(x)$  is the total propensity.

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- It is essentially the BKL algorithm (Bortz-Kalos-Lebowitz) or KMC in condensed matter physics.

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When the reaction rate is very large for a reversible reaction, the reactions will fire back and forth very frequently, but cause very little change of the state.

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How to accelerate the simulation process?

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"Leap Condition: Require the leap time  $\tau$  to be small enough that the change in the state during  $[t, t + \tau)$  will be so slight that no propensity function will suffer an appreciable (i.e., macroscopically noninfinitesimal) change in its value."

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- So we have the tau-leaping scheme

$$\boldsymbol{X}_{t+\delta t} = \boldsymbol{X}_t + \sum_{j=1}^M \boldsymbol{\nu}_j \mathcal{P}(a_j(\boldsymbol{X}_t)\delta t)$$

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► The Taylor expansion of the LHS gives  $\tau | \boldsymbol{\xi} \cdot \nabla a_j | \leq \epsilon a_0(\boldsymbol{X})$ , which gives one stepsize selection strategy

$$\tau \approx \min_{j=1,\dots,M} \left\{ \frac{\epsilon a_0(\boldsymbol{X})}{|\boldsymbol{\xi} \cdot \nabla a_j|} \right\}.$$

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# Tau-leaping Algorithm: Stepsize Selection

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Many more robust stepsize selection strategies are also proposed.

# Some Remarks

#### Remark

The total propensity is  $a_0(\mathbf{X})$ . So the expected waiting time for one reaction fires is  $\mathcal{O}(1/a_0(\mathbf{X}))$ . If

$$\tau \le m/a_0(\boldsymbol{X}), \quad m \sim \mathcal{O}(1)$$

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#### Remark

Compare with the forward Euler step  $x_{n+1} = x_n + f(x_n)\delta t$  for ODE

$$\dot{x} = f(x).$$

We actually fix f(x) as a constant  $f(x_n)$  in  $[t_n, t_{n+1})$  with a similar idea. One will find more connections along this direction.

# Tau-leaping Algorithm: Multiscale Picture

#### Multi-scale picture:

The multi-scale picture from tau-leaping is charming.

From tau-leaping to Chemical Langevin Equation: When  $a_j(\mathbf{X}_t)\tau \gg 1$ ,  $P(a_j(\mathbf{X}_t)\tau) \approx N(a_j(\mathbf{X}_t)\tau, a_j(\mathbf{X}_t)\tau)$  by Central Limit Theorem

$$\boldsymbol{X}_{t+\tau} \approx \boldsymbol{X}_t + \sum_{j=1}^M \boldsymbol{\nu}_j a_j(\boldsymbol{X}_t) \tau + \sum_{j=1}^M \boldsymbol{\nu}_j \sqrt{a_j(\boldsymbol{X}_t) \tau} N(0, 1)$$

which corresponds to CLE

$$d\boldsymbol{X}_t = \sum_{j=1}^M \boldsymbol{\nu}_j a_j(\boldsymbol{X}_t) dt + \sum_{j=1}^M \boldsymbol{\nu}_j \sqrt{a_j(\boldsymbol{X}_t)} d\boldsymbol{W}_t$$

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Tau-leaping Algorithm: Multiscale Picture

From Chemical Langevin Equation to Reaction Rate Equation: When  $a_j(\mathbf{X}_t)\tau \to +\infty$ ,  $N(a_j(\mathbf{X}_t)\tau, a_j(\mathbf{X}_t)\tau) \approx a_j(\mathbf{X}_t)\tau$  by Law of Large Numbers

$$\boldsymbol{X}_{t+\tau} \approx \boldsymbol{X}_t + \sum_{j=1}^M \boldsymbol{\nu}_j a_j(\boldsymbol{X}_t) \tau$$

which corresponds to RRE

$$\frac{d\boldsymbol{X}_t}{dt} = \sum_{j=1}^M \boldsymbol{\nu}_j a_j(\boldsymbol{X}_t)$$

Tau-leaping bridges all of the equations in different scales with a seamless way!

The comparison with fluid mechanics (upscaling) will be instructive.

SSA	$\longrightarrow$	Molecular dynamics
$\downarrow$		$\downarrow$
CLE	$\longrightarrow$	Kinetic theory
$\downarrow$		$\downarrow$
RRE	$\longrightarrow$	Continuum mechanics

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# **Negative Populations**

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## **Negative Populations**

- Because of unboundedness of Poisson R.V., negative populations may appear.
- One choice to avoid N.P. is by binomial tau-leaping.
- Note that Poisson distribution may be viewed as a limit of binomial distribution B(n, p) when n → ∞ with λ = np fixed. That is

$$B(k; n, p) = C_n^k p^k (1-p)^{n-k} \to \mathcal{P}(\lambda).$$

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Algorithm (Binomial Tau-leaping) *Avoiding negative populations.* 



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Avoiding negative populations.

Step 1: At time t, set 
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Firing j-th reaction k<sub>j</sub> times:

$$\tilde{X} + \boldsymbol{\nu}_j k_j \to \tilde{X}$$

Step 3: The iteration is repeated until the final time T is achieved.

Simple Analysis of Binomial Tau-leaping

A simple analysis of binomial tau-leaping as follows. Note that the number of *j*-th reactions:

Poisson : 
$$a_j \tau \pm \sqrt{a_j \tau}$$
  
Binomial :  $a_j \tau \pm \sqrt{a_j \tau} \left(1 - a_j \tau / k_{\max}^{(j)}\right)^{\frac{1}{2}}$ 

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In the law of rare events limit (a<sub>j</sub>τ ≪ k<sup>(j)</sup><sub>max</sub>), they give same result; in the finite size case, the noise is different!

Chemical reactions are usually stiff.



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## Example (Reversible reaction)

Consider the reversible reaction system

$$S_1 \longrightarrow S_2$$
 with rate  $C_1$  (1)

$$S_2 \longrightarrow S_1$$
 with rate  $C_2$  (2)

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when  $C_1$  and  $C_2$  are both large.

▶ Define  $C_1 + C_2 = \lambda$ ,  $X_t^1 + X_t^2 = X^T$  (total number). Here  $\lambda \gg 1$ . Then we have

$$\mathbb{E}X_t^1 = \frac{C_2 X^T}{\lambda} (1 - e^{-\lambda t}) + e^{-\lambda t} X_0^1$$
$$\mathbb{E}X_t^2 = \frac{C_1 X^T}{\lambda} (1 - e^{-\lambda t}) + e^{-\lambda t} X_0^2$$

#### Example (Fast decaying)

Consider the following system

$$S \longrightarrow \emptyset$$
 with rate  $C_1$  (3)

$$\emptyset \longrightarrow S$$
 with rate  $C_2$  (4)

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when  $C_1 - C_2$  is large. Define  $\lambda = C_1 - C_2$ , we have  $\mathbb{E}X_t = e^{-\lambda t}X_0$ .

# Stability Analysis

Now we perform the stability analysis for the stiff reversible reaction system. Suppose the explicit tau-leaping is applied.

$$\begin{cases} X_{n+1}^1 &= X_n^1 - \mathcal{P}(C_1 X_n^1 \delta t) + \mathcal{P}(C_2 X_n^2 \delta t) \\ X_{n+1}^2 &= X_n^2 - \mathcal{P}(C_2 X_n^2 \delta t) + \mathcal{P}(C_1 X_n^1 \delta t) \end{cases}$$

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▶ We have  $X_n^1 + X_n^2 = X^T = Const.$  So we have

$$X_{n+1}^{1} = X_{n}^{1} - \mathcal{P}(C_{1}X_{n}^{1}\delta t) + \mathcal{P}(C_{2}(X^{T} - X_{n}^{1})\delta t).$$

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Taking expectation we obtain

$$\mathbb{E}X_{n+1}^1 = (1 - \lambda \delta t)\mathbb{E}X_n^1 + C_2 X^T \delta t.$$

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The stability condition is

$$|1 - \lambda \delta t| \le 1 \Longrightarrow \delta t \le \frac{2}{\lambda}.$$

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When λ ≫ 1, we have δt ≪ 1. That is the stiffness!
As n → ∞, we have

$$\mathbb{E}X_n^1 \to \frac{C_2}{\lambda} X^T,$$

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which is the correct limit state.

Now consider the variance. At first we have

 $\operatorname{Var}(Y) = \mathbb{E}(\operatorname{Var}(Y|X)) + \operatorname{Var}(\mathbb{E}(Y|X)).$ 

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$$\begin{aligned} \operatorname{Var}(X_{n+1}^{1}) &= C_{1}\delta t \mathbb{E}X_{n}^{1} + C_{2}\delta t \mathbb{E}(X^{T} - X_{n}^{1}) \\ &+ \operatorname{Var}\left(X_{n}^{1} - C_{1}X_{n}^{1}\delta t + C_{2}(X^{T} - X_{n}^{1})\delta t\right) \\ &= (1 - \lambda \delta t)^{2} \operatorname{Var}(X_{n}^{1}) + (C_{1} - C_{2})\delta t \mathbb{E}X_{n}^{1} + C_{2}\delta t X^{T}. \end{aligned}$$

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▶ As  $n \to \infty$ , we have

$$\operatorname{Var}(X_n^1) \to \frac{2}{2 - \lambda \delta t} \frac{C_1 C_2 X^T}{(C_1 + C_2)^2} = \frac{2}{2 - \lambda \delta t} \operatorname{Var}(X_\infty^1) \ge \operatorname{Var}(X_\infty^1).$$

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▶ In order to get the right variance, we need  $\lambda \delta t \rightarrow 0$ . Since  $\lambda \gg 1$ , we need  $\delta t \rightarrow 0$ , which is a strict constraint.

# Stiff System

Strategy: Implicit method to overcome stiffness.

The first choice is

$$X_{n+1}^{1} = X_{n}^{1} - \mathcal{P}(C_{1}X_{n+1}^{1}\delta t) + \mathcal{P}(C_{2}(X^{T} - X_{n+1}^{1})\delta t).$$

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• But the problem is how to sample  $\mathcal{P}(C_1 X_{n+1}^1 \delta t)$ . If we apply the iteration

$$X_{n+1}^{1,k+1} = X_n^1 - \mathcal{P}(C_1 X_{n+1}^{1,k} \delta t) + \mathcal{P}(C_2 (X^T - X_{n+1}^{1,k}) \delta t),$$

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The second choice is semi-implicit method as

$$\begin{aligned} X_{n+1}^{1} &= X_{n}^{1} - C_{1} X_{n+1}^{1} \delta t + C_{2} (X^{T} - X_{n+1}^{1}) \delta t \\ &- \left[ \mathcal{P}(C_{1} X_{n}^{1} \delta t) - C_{1} X_{n}^{1} \delta t \right] \\ &+ \left[ \mathcal{P}(C_{2} (X^{T} - X_{n}^{1}) \delta t) - C_{2} (X^{T} - X_{n}^{1}) \delta t \right]. \end{aligned}$$

# Stiff System

Similar analysis as before shows the stability condition

$$\left|\frac{1}{1+\lambda t}\right| \le 1.$$

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Trapezoidal method is a good choice for linear problem. But the story goes on for nonlinear stiff problem!

### **Convergence** Analysis

Issue: the chemical reaction kinetics is a pure jump process with state dependent intensity.

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- Construct jump process with state dependent intensity from constant jump intensity process (P. Protter, 1983) — Acceptance rejection method

$$\mu(dt) = \int_0^A \mathbb{1}_{\{0 < x \le a_0(\boldsymbol{X}_t)\}} \lambda(dt \times dx).$$

 $\lambda(dt \times dx)$  is the Poisson random measure generated from a constant jump intensity process.  $\mu(dt)$  has intensity  $a_0(\mathbf{X}_t)$ .

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λ(dt × dx) is the Poisson random measure generated from a constant jump intensity process. μ(dt) has intensity a<sub>0</sub>(X<sub>t</sub>).
The SDE form for the CME

$$d\boldsymbol{X}_t = \sum_{j=1}^M \int_0^A \boldsymbol{\nu}_j c_j(x; \boldsymbol{X}_{t-}) \lambda(dt \times dx),$$

where

$$c_j(x; \boldsymbol{X}_t) = \begin{cases} 1, & \text{if } x \in (\sum_{i=1}^{j-1} a_i(\boldsymbol{X}_t), \sum_{i=1}^j a_i(\boldsymbol{X}_t)], \\ 0, & \text{otherwise.} \end{cases}$$

# Tau-leaping is an explicit Euler scheme

Decomposition

$$d\boldsymbol{X}_{t} = \sum_{j=1}^{M} \int_{0}^{A} \boldsymbol{\nu}_{j} c_{j}(x; \boldsymbol{X}_{t-}) m(dt \times dx)$$
  
+ 
$$\sum_{j=1}^{M} \int_{0}^{A} \boldsymbol{\nu}_{j} c_{j}(x; \boldsymbol{X}_{t-}) (\lambda - m)(dt \times dx)$$
  
= 
$$\boldsymbol{P}_{1} + \boldsymbol{P}_{2}.$$

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- We call  $P_1$  the drift term and  $P_2$  is the jump term.
- Explicit Euler scheme tau-leaping method!

$$\boldsymbol{X}_{n+1} = \boldsymbol{X}_n + \sum_{j=1}^M \boldsymbol{\nu}_j \mathcal{P}(a_j(\boldsymbol{X}_n)\delta t_n)$$

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### Other tau-leaping schemes

Implicit tau-leaping: semi-implicit Euler of SDEs

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Implicit tau-leaping: semi-implicit Euler of SDEs

Stochastic theta methods:

$$\begin{aligned} \boldsymbol{X}_{n+1} &= \boldsymbol{X}_n + \sum_{j=1}^M \theta \boldsymbol{\nu}_j \Big( a_j(\boldsymbol{X}_{n+1}) - a_j(\boldsymbol{X}_n) \Big) \delta t_n \\ &+ \sum_{j=1}^M \boldsymbol{\nu}_j \mathcal{P}(a_j(\boldsymbol{X}_n) \delta t_n). \end{aligned}$$

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Milstein scheme: Not directly imply any implementable scheme!

# **Convergence** Theorem

Theorem (Mean square convergence) With the assumptions before we have

$$\sup_{n \le N_T} \mathbb{E} |\boldsymbol{X}_n - \boldsymbol{X}_{t_n}|^2 \le C\tau,$$

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#### Theorem (Weak convergence)

Under the assumptions, for any continuous function g(x) satisfying exponential growth condition

$$|g(\boldsymbol{x})| \leq C_g B^{|\boldsymbol{x}|}, \ \ \boldsymbol{x} \in \mathbb{R}^N$$
 and  $C_g, B > 0.$ 

We have

$$\left|\mathbb{E}g(\boldsymbol{X}_{N_T}) - \mathbb{E}g(\boldsymbol{X}_T)\right| \leq C\tau,$$

where  $T = t_{N_T}$ ,  $\tau = \max_n \delta t_n$ .

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The chemical master equation (FPE) is as before. Denote it as  $\partial_t P = \mathcal{L}P$ . Here  $\mathcal{L}$  is the adjoint operator of the infinitesimal generator

$$\mathcal{L}^* u = \sum_{j=1}^M a_j(\boldsymbol{x}) u(\boldsymbol{x} + \boldsymbol{\nu}_j, t) - \sum_{j=1}^M a_j(\boldsymbol{x}) u(\boldsymbol{x}, t)$$
$$= \sum_{j=1}^M a_j(\boldsymbol{x}) \Big( u(\boldsymbol{x} + \boldsymbol{\nu}_j, t) - u(\boldsymbol{x}, t) \Big).$$

For the stationary solution, we ask

$$\mathcal{L}P=0.$$

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► For reversible reaction, we only consider the equation for x since x + y = x<sup>T</sup> (v<sub>1</sub> = −1, v<sub>2</sub> = 1):

$$\left(C_1(x+1)p(x+1) - C_1xp(x)\right) + \left(C_2(x^T - x + 1)p(x-1) - C_2(x^T - x)p(x)\right) = 0.$$

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▶ If x = 0,  $a_1(x)p(x) - a_2(x-1)p(x-1) = a(0)p(0) = 0$ , we have the detailed balance  $a_1(x)p(x) = a_2(x-1)p(x-1)$ , so

$$\frac{p(x)}{p(x-1)} = \frac{a_2(x-1)}{a_1(x)} \Longrightarrow \frac{p(x)}{p(0)} = \frac{a_2(x-1)}{a_1(x)} \frac{a_2(x-2)}{a_1(x-1)} \cdots \frac{a_2(0)}{a_1(1)}$$

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We obtain the stationary distribution

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From the normalization we have

$$p(x) \sim B(x^T, q), \quad q = \frac{C_2}{C_1 + C_2}$$

with mean  $x^T C_2/(C_1+C_2)$ , and variance  $x^T C_1 C_2/(C_1+C_2)$ .