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## FULLY DISCRETE CONVERGENCE ESTIMATES FOR VORTEX METHODS IN BOUNDED DOMAINS\*

YING LUNG-AN<sup>†</sup> AND ZHANG PING-WEN<sup>†</sup>

**Abstract.** In this paper the authors study vortex method for 2-dimensional Euler equations of incompressible flow in bounded domains. To approximate the initial vortex field by a sum of vortex blobs with arbitrary high accuracy, this field is extended smoothly to a small neighborhood of the boundary. And the computation is carried out in the extended domain. To construct the velocity field from vorticity, a second-order isoparametric finite element method is applied, and to solve the characteristic equations, the explicit Euler's scheme is considered. Optimal error bounds for this fully discrete scheme are obtained.

**Key words.** Euler equation, vortex method, convergence, initial-boundary value problem, finite element method, Euler's scheme

**AMS subject classifications.** 65N15, 35Q35, 76M25, 76C05

**1. Introduction.** Vortex methods are efficient numerical techniques for simulating incompressible flow, especially flow with high Reynold's number. The mathematical foundation of the vortex methods has been studied by many authors. By virtue of a viscous splitting approach, the equations are decomposed into Euler equations and pure diffusion equations. Thus the study of vortex methods for Euler equations is an important part in the theory, and the most fruitful results have been obtained in this direction.

The convergence of vortex methods for the initial value problems of Euler equations was first obtained by Hald [6]. Then the results were improved and different proofs were given by several authors [2], [3], [4], [7]. Recently, the first author of this paper considered the convergence problem for two-dimensional bounded domains [10] and obtained optimal error estimates for a semidiscretization scheme, where it was assumed that the equations for stream functions and the system of ordinary differential equations for partial trajectories were solved exactly. In [10] a finite element approximation for the equations for the stream function was also considered and convergence results were given, but constants in the error bounds depended on the vortex blob parameters.

One purpose of this paper is to prove that the rate of convergence of the finite element schemes can be independent of the vortex blob parameters, provided second-order isoparametric finite elements are used instead of linear ones. The other purpose is to give error estimates for fully discretized two-dimensional vortex methods for initial-boundary value problems of Euler equations.

The paper is organized as follows. In §2, we recall a result in [10]. In §3, we prove an error estimate for the combined effect of vortex discretization and finite element approximation. In §4 we prove the error estimate for full discretization problems.

**2. A convergence theorem for semidiscretization.** Let  $\Omega \subset \mathbb{R}^2$  be a convex and bounded domain, whose boundary  $\partial\Omega$  is sufficiently smooth. Denote by  $x =$

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$(x_1, x_2)$  the points in  $\mathbb{R}^2$ . We consider the following initial-boundary value problems

$$(2.1) \quad \frac{\partial u}{\partial t} + (u \cdot \nabla)u + \frac{1}{\rho} \nabla \pi = f,$$

$$(2.2) \quad \nabla \cdot u = 0,$$

$$(2.3) \quad u \cdot n \big|_{x \in \partial\Omega} = 0,$$

$$(2.4) \quad u \big|_{t=0} = u_0,$$

where  $u = (u_1, u_2)$  stands for velocity,  $\pi$  stands for pressure,  $f = (f_1, f_2)$  is the external force, the density  $\rho$  is a positive constant,  $n$  is the unit outward normal vector along  $\partial\Omega$ , and the initial data  $u_0$  satisfies

$$\nabla \cdot u_0 = 0, u_0 \cdot n \big|_{\partial\Omega} = 0.$$

If  $f, u_0$  are sufficiently smooth, then the solutions  $u$  and  $\pi$  are also sufficiently smooth on the domain  $\bar{\Omega} \times [0, T]$ , where  $T$  is an arbitrary positive constant.

Let  $\omega = -\nabla \wedge u, \omega_0 = -\nabla \wedge u_0$ , and  $\psi$  be the stream function corresponding to  $u$ . Then (2.1)–(2.4) is equivalent to

$$(2.5) \quad \frac{\partial \omega}{\partial t} + u \cdot \nabla \omega = -\nabla \wedge f \equiv F,$$

$$(2.6) \quad -\Delta \psi = \omega, u = \nabla \wedge \psi,$$

$$(2.7) \quad \psi \big|_{x \in \partial\Omega} = 0,$$

$$(2.8) \quad \omega \big|_{t=0} = \omega_0.$$

We extend functions  $u_0$  and  $f$ , still denoted by  $u_0$  and  $f$ , such that they are sufficiently smooth on  $\mathbb{R}^2$  and  $\mathbb{R}^2 \times [0, T]$ , respectively, and the supports of them are compact. Let  $c$  be any positive constant. We define

$$\Omega_c = \{x, \text{dist}(x, \bar{\Omega}) < c\}.$$

The “blob function” is defined as follows,  $\zeta(x)$  is a cutoff function, such that  $\zeta \equiv 0$  for  $|x| > 1$  and

$$\zeta_\varepsilon(x) = \frac{1}{\varepsilon^2} \zeta\left(\frac{x}{\varepsilon}\right).$$

With that notation, the semidiscretization scheme for (2.5)–(2.8) is

$$(2.9) \quad \omega^\varepsilon(x, t) = \sum_{j \in J_1} \alpha_j^\varepsilon(t) \zeta_\varepsilon(x - X_j^\varepsilon(t)),$$

$$(2.10) \quad \frac{d\alpha_j^\varepsilon}{dt} = h^2 F(X_j^\varepsilon(t), t), \alpha_j^\varepsilon(0) = \alpha_j,$$

$$(2.11) \quad \frac{dX_j^\varepsilon}{dt} = g^\varepsilon(X_j^\varepsilon(t), t), X_j^\varepsilon(0) = X_j,$$

$$(2.12) \quad -\Delta \psi^\varepsilon = \omega^\varepsilon, \psi^\varepsilon \big|_{x \in \partial\Omega} = 0,$$

$$(2.13) \quad u^\varepsilon = \nabla \wedge \psi^\varepsilon,$$

$$(2.14) \quad g^\varepsilon(x, t) = \sum_{i=1}^M a_i u^\varepsilon(x^{(i)}, t),$$

where  $j = (j_1, j_2)$ ,  $X_j = (j_1 h, j_2 h)$ ,  $\alpha_j = h^2 \omega_0(X_j)$ , and  $J_1 = \{j, X_j \in \Omega_d\}$ , if  $x \in \bar{\Omega}$ , then  $x^{(i)} = x$ , otherwise

$$x^{(i)} = (i + 1)Y - ix,$$

where  $Y$  is the nearest point on  $\partial\Omega$  to  $x$ ; the terms  $a_i$  are the solutions of the system

$$\sum_{i=1}^M (-i)^j a_i = 1 \quad j = 0, \dots, M;$$

and  $\varepsilon > 0, h > 0, d > 0$  are mesh parameters. Equations (2.14) makes sense only if  $x^{(i)}$  belongs to  $\bar{\Omega}$ , but it is proved in [10] that this fact is true provided  $d$  is small enough. In this scheme the function  $g^\varepsilon$  plays the role of velocity, which is equal to  $u^\varepsilon$  in the domain and interpolated to the exterior part of  $\Omega$ . This is a natural way to deal with blobs near the boundary. Using  $g^\varepsilon$  and a “slightly larger” domain  $\Omega_d$  in computation, all blobs move according to a uniform formula (2.11).

Now we state the convergence results. The notation  $W^{m,p}(\Omega)$  for conventional Sobolev spaces and  $\|\cdot\|_{m,p}$  for the norms of them are applied throughout this paper. Let  $X_j(t)$  be characteristic curves that satisfy

$$\frac{dX_j(t)}{dt} = u(X_j(t), t), \quad X_j(0) = X_j.$$

As a rule, we admit the value of  $u$  as an extension if  $X_j(t) \in \bar{\Omega}$ . Then set

$$J_2 = \{j; X_j \in \Omega_{C_0\varepsilon} \cap \Omega_d\},$$

$$\|e(t)\|_p = \left( h^2 \sum_{j \in J_2} |X_j(t) - X_j^\varepsilon(t)|^p \right)^{1/p}, \quad 1 \leq p < \infty,$$

where  $C_0$  is a positive constant to be determined. The following theorem is proved in [10].

**THEOREM 1.** *If we have  $m \geq 1, k \geq 2$ , such that  $\zeta \in W^{m+1,\infty}(\mathbb{R}^2)$  and*

$$(2.15) \quad \int_{\mathbb{R}^2} \zeta(x) dx = 1,$$

$$\int_{\mathbb{R}^2} x^\alpha \zeta(x) dx = 0 \quad \forall \alpha \in N^2 \quad \text{with } 1 \leq |\alpha| \leq k - 1$$

and if there is a constant  $\tilde{C}$ , such that

$$(2.16) \quad \tilde{C}^{-1} \varepsilon^a \leq h \leq \tilde{C} \varepsilon^{1 + \frac{k-1}{m}},$$

where  $a \geq 1 + \frac{k-1}{m}$ , and if the constant in expression (2.14),  $M \geq k$ , then for any  $p \in [1, \infty)$ , there are positive constants  $d_0, C_0, C_1$ , and  $C_2$  such that if  $d \leq d_0$ , then the solution of problem (2.9)–(2.14) satisfy

$$(2.17) \quad |\nabla u^\varepsilon(x, t)| \leq C_1, \quad x \in \bar{\Omega},$$

$$(2.18) \quad \|u - u^\varepsilon\|_{0,p,\Omega} + \|e(t)\|_p \leq C_2 \varepsilon^k$$

for  $t \in [0, T]$ .

For our later use, we need the following.

**COROLLARY.** *Under the assumption of Theorem 1, let  $C_3 > 0$  be given. Then there is a constant  $C_4$ , such that*

$$(2.19) \quad \|\omega^\varepsilon(\cdot, t)\|_{k-1,p,\Omega_{C_3\varepsilon}} \leq C_4, \quad t \in [0, T].$$

*Proof.* It was proved in [10] that the points  $X_j$  out of  $\Omega_{C_0\varepsilon}$  do not contribute to the value of  $\omega^\varepsilon$  in  $\Omega$ . Now we replace  $C_0$  by  $C_0 + C_3$ . Then the points  $X_j$  do not contribute to the value of  $\omega^\varepsilon$  in  $\Omega_{C_3\varepsilon}$ .

By the proof of Lemmas 2, 3, and 4 of [10],

$$(2.20) \quad \begin{aligned} & |\omega - \omega^\varepsilon|_{l-1,p,\Omega_{C_3\varepsilon}} \\ & \leq C \left\{ \varepsilon^k + \left(1 + \frac{h}{\varepsilon}\right)^{\frac{2}{r}} \frac{h^m}{\varepsilon^{m+l-1}} + \frac{h^N}{\varepsilon^{N+l-1}} \right. \\ & \quad \left. + \frac{1}{\varepsilon^l} \left(1 + \frac{1}{\varepsilon} \|e(t)\|_{0,\infty,h}\right)^{\frac{2}{q}} \|e(t)\|_p + \frac{1}{\varepsilon^l} \int_0^t \|e(s)\|_p ds \right\}, \end{aligned}$$

where  $r \in [1, 2]$ ,  $\frac{1}{p} + \frac{1}{q} = 1$ ,  $N \geq 3$ ,  $C$  is a positive constant, and

$$\|e(t)\|_\infty = \max_{j \in J_2} |X_j(t) - X_j^\varepsilon(t)|.$$

By (2.31) of [10], we have

$$\|e(t)\|_\infty \leq \frac{1}{h^{\frac{2}{p'}}} \|e(t)\|_{p'}$$

for any  $p' \in [1, \infty)$ . We take  $l = k$ ,  $p' \geq 4$ ,  $N \geq m$ . Then we get the upper bound of the right-hand side of (2.20).  $\square$

**3. Further discretization by finite element methods.** Let  $C$  denote a generic constant independent of mesh parameters. For simplicity we only consider quadratic triangular isoparametric elements of Lagrange type here. Let a triangle  $\hat{K}$  be the reference element, and the set of six nodes consists of three vertices and three midpoints of the edges. Denote by  $F_K$  the isoparametric mapping from  $\hat{K}$  to each element  $K$ . If  $K$  is an interior one, then we take  $F_K$  affine. If  $K$  is a ‘‘boundary’’ element, the nodes of which is shown in Fig. 3.1 where  $\tilde{a}_{12,K}$  is the midpoint of  $\overline{a_{1,K}a_{2,K}}$ , the node  $a_{12,K}$  belongs to  $\partial\Omega$ , and  $\overline{\tilde{a}_{12,K}a_{12,K}}$  is perpendicular to  $\overline{a_{1,K}a_{2,K}}$ . The nodes  $a_{13,K}$  and  $a_{23,K}$  are simply midpoints. It is known that  $F_K$  is uniquely determined by the nodes.

Let  $\mathcal{T}_\delta = \{K\}$  and  $\overline{\Omega}^\delta = \bigcup_{K \in \mathcal{T}_\delta} K$ , where  $\delta$  is the size of the largest diameter of elements. We assume that the partition is regular and quasiuniform. Then we define the finite element space

$$V^\delta = \{v \in H_0^1(\Omega^\delta); v|_K \in P_2(\hat{K}) \circ F_K^{-1}\},$$

where  $P_2(\hat{K})$  is the space of all polynomials of degree  $\leq 2$  and  $\Omega^\delta$  is the interior of  $\overline{\Omega}^\delta$ . We consider in this section the following scheme:

$$(3.1) \quad \omega^\delta(x, t) = \sum_{j \in J_1} \alpha_j^\delta(t) \zeta_\varepsilon(x - X_j^\delta(t)),$$

$$(3.2) \quad \frac{d\alpha_j^\delta(t)}{dt} = h^2 F(X_j^\delta(t), t), \quad \alpha_j^\delta(0) = \alpha_j,$$

$$(3.3) \quad \frac{dX_j^\delta(t)}{dt} = g^\delta(X_j^\delta(t), t), \quad X_j^\delta(0) = X_j,$$

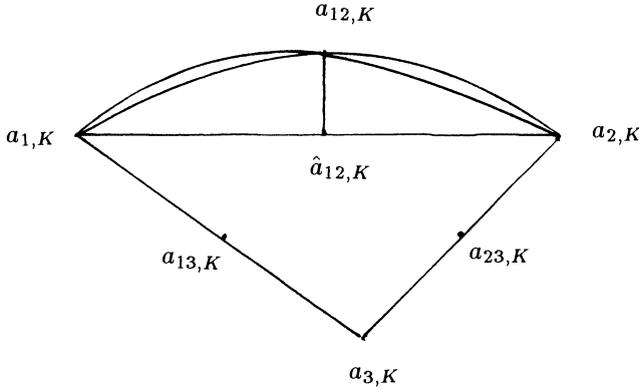


FIG. 3.1

$\psi^\delta \in V^\delta$ , and

$$(3.4) \quad \int_{\Omega^\delta} \nabla \psi^\delta \cdot \nabla v dx = \int_{\Omega^\delta} \omega^\delta v dx \quad \forall v \in V^\delta,$$

$$(3.5) \quad u^\delta = \nabla \wedge \psi^\delta,$$

where

$$g^\delta(x, t) = \sum_{i=1}^M a_i u^\delta(x_\delta^{(i)}, t).$$

As before  $x_\delta^{(i)} = x$  if  $x \in \bar{\Omega}^\delta$ , and otherwise

$$x_\delta^{(i)} = (i + 1)Y_\delta - ix,$$

where  $Y_\delta$  is the nearest point on  $\partial\Omega^\delta$  to  $x$ .  $Y_\delta$  is not necessary unique; we pick up an arbitrary one. Now  $g^\delta$  is no longer continuous, so (3.3) is satisfied in a generalized sense; see [10]. Here (3.4) and (3.5) are just the definition of a weak finite element solution to (2.12). Since  $\Omega$  is arbitrary, spacial discretization is needed in solving (2.12). We will prove convergence and get error estimates for it.

Since  $\Omega^\delta \not\subset \Omega$  we need to extend some functions from  $\Omega$  to the whole space  $\mathbb{R}^2$ . Since  $\partial\Omega$  is sufficiently smooth, there exists a strong  $m$ -extension operator  $E$  on  $\Omega$ , such that [1]

$$(3.6) \quad \|E\psi\|_{k,p,\mathbb{R}^2} \leq C\|\psi\|_{k,p,\Omega}, \quad \forall 0 \leq k \leq m, \quad 1 \leq p < \infty, \quad \psi \in W^{m,p}(\Omega).$$

In this section, we take  $m$  large enough, and extend the stream function  $\psi^\varepsilon$ , still denoted by  $\psi^\varepsilon$ , then set  $u^\varepsilon$  and  $\omega^\varepsilon$  to be the corresponding velocity and vorticity, all of them now defined on  $\mathbb{R}^2 \times [0, T]$ .

We are now in a position to estimate  $u^\varepsilon - u^\delta$  and  $X_j^\varepsilon(t) - X_j^\delta(t)$ . We define

$$\|e(t)\|_p = \left( h^2 \sum_{j \in J_1} |X_j^\varepsilon(t) - X_j^\delta(t)|^p \right)^{\frac{1}{p}},$$

$$\|e(t)\|_\infty = \max_{j \in J_1} |X_j^\varepsilon(t) - X_j^\delta(t)|.$$

LEMMA 1. *Under the assumptions of Theorem 1,*

$$(3.7) \quad \begin{aligned} \|\psi^\varepsilon(\cdot, t) - \psi^\delta(\cdot, t)\|_{1,p,\Omega^\delta} &\leq C\delta^2 + C \left\{ \left(1 + \frac{1}{\varepsilon} \|e(t)\|_\infty\right)^{\frac{2}{q}} \|e(t)\|_p \right. \\ &\quad \left. + \int_0^t \|e(s)\|_p ds \right\}, \end{aligned}$$

provided  $\delta \leq C_3\varepsilon$ , where  $p \geq 2, \frac{1}{p} + \frac{1}{q} = 1$ .

*Proof.* We define two operators: If  $\psi^\varepsilon$  corresponds to  $\omega^\varepsilon$  on the basis of (2.12), then we define  $\psi^\varepsilon = \Delta^{-1}\omega^\varepsilon$ , likewise (3.4) defines an operator  $\psi^\delta = \Delta_\delta^{-1}\omega^\delta$ . Then

$$\begin{aligned} \psi^\varepsilon - \psi^\delta &= \Delta^{-1} \left( \sum_{j \in J_1} \alpha_j^\varepsilon(t) \zeta_\varepsilon(\cdot - X_j^\varepsilon(t)) \right) - \Delta_\delta^{-1} \left( \sum_{j \in J_1} \alpha_j^\delta(t) \zeta_\varepsilon(\cdot - X_j^\delta(t)) \right) \\ &\equiv \phi_1 + \phi_2 + \phi_3, \end{aligned}$$

on the domain  $\Omega^\delta$ , where

$$\begin{aligned} \phi_1 &= \Delta^{-1} \left( \sum_{j \in J_1} \alpha_j^\delta(t) (\zeta_\varepsilon(\cdot - X_j^\varepsilon(t)) - \zeta_\varepsilon(\cdot - X_j^\delta(t))) \right), \\ \phi_2 &= \Delta^{-1} \left( \sum_{j \in J_1} (\alpha_j^\varepsilon(t) - \alpha_j^\delta(t)) \zeta_\varepsilon(\cdot - X_j^\delta(t)) \right), \\ \phi_3 &= (\Delta^{-1} - \Delta_\delta^{-1}) \sum_{j \in J_1} \alpha_j^\delta(t) \zeta_\varepsilon(\cdot - X_j^\delta(t)). \end{aligned}$$

Proceeding in a manner similar to [10], we have

$$(3.8) \quad \|\phi_1(\cdot, t)\|_{1,p,\Omega} \leq C \left(1 + \frac{1}{\varepsilon} \|e(t)\|_\infty\right)^{\frac{2}{q}} \|e(t)\|_p,$$

$$(3.9) \quad \|\phi_2(\cdot, t)\|_{1,p,\Omega} \leq C \int_0^t \|e(s)\|_p ds,$$

and

$$(3.10) \quad \begin{aligned} \|\omega^\varepsilon(\cdot, t) - \omega^\delta(\cdot, t)\|_{1,p,\Omega_{C_3\varepsilon}} &\leq \frac{C}{\varepsilon^2} \left\{ \left(1 + \frac{1}{\varepsilon} \|e(t)\|_\infty\right)^{\frac{2}{q}} \|e(t)\|_p \right. \\ &\quad \left. + \int_0^t \|e(s)\|_p ds \right\}. \end{aligned}$$

Inequality (3.10) and the Corollary of Theorem 1 yield

$$(3.11) \quad \begin{aligned} \|\omega^\delta(\cdot, t)\|_{1,p,\Omega_{C_3\varepsilon}} &\leq C + \frac{C}{\varepsilon^2} \left\{ \left(1 + \frac{1}{\varepsilon} \|e(t)\|_\infty\right)^{\frac{2}{q}} \|e(t)\|_p \right. \\ &\quad \left. + \int_0^t \|e(s)\|_p ds \right\}. \end{aligned}$$

To estimate  $\phi_3$ , we define a function  $\psi_1$  that solves

$$-\Delta\psi_1 = \omega^\delta, \quad x \in \Omega, \quad \psi_1|_{x \in \partial\Omega} = 0.$$

Then function  $\psi^\delta$  determined by (3.4) is the finite element approximation of  $\psi_1$ . An abstract error estimate of the isoparametric finite element method shows (see [5, Thm. 4.4.1])

$$(3.12) \quad \|\psi_1 - \psi^\delta\|_{1,\Omega^\delta} \leq C \left( \inf_{v \in V^\delta} \|\psi_1 - v\|_{1,\Omega^\delta} + \sup_{\substack{v \in V^\delta \\ v=0}} \frac{|\int_{\Omega^\delta} \nabla \psi_1 \cdot \nabla v dx - \int_{\Omega^\delta} \omega^\delta v dx|}{\|v\|_{1,\Omega^\delta}} \right).$$

Integrating by parts and using the Hölder inequality gives

$$\begin{aligned} & \left| \int_{\Omega^\delta} \nabla \psi_1 \cdot \nabla v dx - \int_{\Omega^\delta} \omega^\delta v dx \right| \\ &= \left| - \int_{\Omega^\delta} (\Delta \psi_1 + \omega^\delta) v dx \right| \\ &= \left| \int_{\Omega^\delta \setminus \Omega} (\Delta \psi_1 + \omega^\delta) v dx \right| \\ &\leq \|\Delta \psi_1 + \omega^\delta\|_{0,6,\Omega^\delta \setminus \Omega} \|v\|_{0,6,\Omega^\delta \setminus \Omega} \cdot (\text{meas}(\Omega^\delta \setminus \Omega))^{\frac{2}{3}}. \end{aligned}$$

It will be shown in the Appendix that

$$\text{meas}(\Omega^\delta \setminus \Omega) \leq C\delta^3.$$

This fact and the conclusion of the embedding theorem,  $H^1(\Omega^\delta) \rightarrow L^6(\Omega^\delta)$ , shows that the second term of the right-hand side of (3.12) is bounded by

$$C\delta^2(\|\omega^\delta\|_{1,\Omega^\delta} + \|\psi_1\|_{3,\Omega^\delta}) \leq C\delta^2\|\omega^\delta\|_{1,\Omega^\delta \cup \Omega}.$$

An interpolation error estimates theorem [5, Thm. 4.3.4] gives the same bound for the first term; therefore,

$$\|\psi_1 - \psi^\delta\|_{1,\Omega^\delta} \leq C\delta^2\|\omega^\delta\|_{1,\Omega_{C_3\varepsilon}}.$$

By the  $L^\infty$  estimate [8], we have

$$\|\psi_1 - \psi^\delta\|_{1,\infty,\Omega^\delta} \leq C\delta^2\|\omega^\delta\|_{1,\infty,\Omega_{C_3\varepsilon}}.$$

Then the Stampachia interpolation inequality gives

$$\|\psi_1 - \psi^\delta\|_{1,p,\Omega^\delta} \leq C\delta^2\|\omega^\delta\|_{1,p,\Omega_{C_3\varepsilon}},$$

that is,

$$(3.13) \quad \|\phi_3\|_{1,p,\Omega^\delta} \leq C\delta^2\|\omega^\delta\|_{1,p,\Omega_{C_3\varepsilon}}.$$

Finally (3.8), (3.9), (3.13), and (3.11) yield (3.7).  $\square$

LEMMA 2. *Under the assumptions of Theorem 1 with  $k \geq 3$ ,*

$$(3.14) \quad \|e(t)\|_p \leq C\delta^2 + C \int_0^t \left(1 + \frac{1}{\delta} \|e(s)\|_\infty + \frac{h}{\delta}\right)^{\frac{2}{p}} \cdot \left\{ \delta^2 + \left(1 + \frac{1}{\varepsilon} \|e(s)\|_\infty\right)^{\frac{2}{q}} \|e(s)\|_p + \int_0^s \|e(t)\|_p dt \right\} ds,$$

provided  $\delta \leq C_3\varepsilon$ , and  $d$  is small enough, where  $p \geq 2$ ,  $\frac{1}{p} + \frac{1}{q} = 1$ .

*Proof.* We define

$$g^{\varepsilon\delta}(x, t) = \sum_{i=1}^M a_i u^\varepsilon(x_\delta^{(i)}, t).$$

Then taking (2.11) and (3.3) into account, we have

$$\frac{dX_j^\varepsilon(t)}{dt} - \frac{dX_j^\delta(t)}{dt} = I_1 + I_2, \quad X_j^\varepsilon(0) - X_j^\delta(0) = 0,$$

where

$$\begin{aligned} I_1 &= g^\varepsilon(X_j^\varepsilon(t), t) - g^{\varepsilon\delta}(X_j^\delta(t), t), \\ I_2 &= g^{\varepsilon\delta}(X_j^\varepsilon(t), t) - g^\delta(X_j^\delta(t), t). \end{aligned}$$

By (2.17) and the definitions of functions  $g^\varepsilon$  and  $g^{\varepsilon\delta}$ , we obtain

$$\begin{aligned} |I_1| &\leq C|(X_j^\varepsilon(t))^{(i)} - (X_j^\delta(t))_\delta^{(i)}| \\ &\leq C|(X_j^\varepsilon(t))^{(i)} - (X_j^\delta(t))^{(i)}| + C|(X_j^\delta(t))^{(i)} - (X_j^\delta(t))_\delta^{(i)}|. \end{aligned}$$

Since  $\Omega$  is convex, the first term is bounded by  $C|X_j^\varepsilon(t) - X_j^\delta(t)|$ , and we will prove in the Appendix that the second term is bounded by  $C\delta^2$ ; therefore,

$$(3.15) \quad |I_1| \leq C|X_j^\varepsilon(t) - X_j^\delta(t)| + C\delta^2.$$

Now we estimate  $I_2$ , denoted by  $K_j^{(i)}$ , the element to which the point  $(X_j^\delta(t))_\delta^{(i)}$  belongs. Then

$$\begin{aligned} |I_2| &\leq C \sum_{i=1}^M \|u^\varepsilon(\cdot, t) - u^\delta(\cdot, t)\|_{0, \infty, K_j^{(i)}} \\ &= C \sum_{i=1}^M |\psi^\varepsilon(\cdot, t) - \psi^\delta(\cdot, t)|_{1, \infty, K_j^{(i)}}. \end{aligned}$$

Let  $I: C(K) \rightarrow P_2(\hat{K}) \circ F_K^{-1}$  be the interpolation operator associated with the nodes. We have

$$|\psi^\varepsilon(\cdot, t) - \psi^\delta(\cdot, t)|_{1, \infty, K_j^{(i)}} \leq |\psi^\varepsilon(\cdot, t) - I\psi^\varepsilon(\cdot, t)|_{1, \infty, K_j^{(i)}} + |I\psi^\varepsilon(\cdot, t) - \psi^\delta(\cdot, t)|_{1, \infty, K_j^{(i)}}.$$

By the interpolation estimate for isoparametric elements and the Corollary of Theorem 1, we obtain

$$\begin{aligned} |\psi^\varepsilon(\cdot, t) - I\psi^\varepsilon(\cdot, t)|_{1, \infty, K_j^{(i)}} &\leq C\delta^2 |\psi^\varepsilon(\cdot, t)|_{3, \infty, K_j^{(i)}} \\ &\leq C\delta^2 |\psi^\varepsilon(\cdot, t)|_{3, \infty, \Omega} \\ &\leq C\delta^2. \end{aligned}$$

Using the inverse inequality gives

$$|I\psi^\varepsilon(\cdot, t) - \psi^\delta(\cdot, t)|_{1, \infty, K_j^{(i)}} \leq C\delta^{-\frac{2}{p}} \|I\psi^\varepsilon(\cdot, t) - \psi^\delta(\cdot, t)\|_{1, p, K_j^{(i)}}.$$

Therefore

$$|I_2| \leq C\delta^2 + \sum_{i=1}^M \delta^{-\frac{2}{p}} \|I\psi^\varepsilon(\cdot, t) - \psi^\delta(\cdot, t)\|_{1, p, K_j^{(i)}}.$$

In conjunction with (3.15), this gives

$$\begin{aligned} (3.16) \quad \left| \frac{dX_j^\varepsilon(t)}{dt} - \frac{dX_j^\delta(t)}{dt} \right| &\leq C\delta^2 + C|X_j^\varepsilon(t) - X_j^\delta(t)| \\ &\quad + C \sum_{i=1}^M \delta^{-\frac{2}{p}} \|I\psi^\varepsilon(\cdot, t) - \psi^\delta(\cdot, t)\|_{1, p, K_j^{(i)}}. \end{aligned}$$

Before summing up (3.16) with respect to  $j$ , we should estimate the number of points  $(X_j^\delta(t))_\delta^{(i)}$  which lie in one single element  $K$ . First of all, let us estimate  $\text{card}\{j \in J_1; X_j^\delta(t) \in K\}$ . We define

$$B_j = \left\{ x \in \mathbb{R}^2; \left( j_i - \frac{1}{2} \right) h \leq x_i \leq \left( j_i + \frac{1}{2} \right) h, i = 1, 2 \right\}.$$

Consider the initial value problem

$$\frac{dy}{dt} = g^\varepsilon(y, t), \quad y|_{t=0} = y_0 \in B_j$$

and define  $B_j(t) = \{y(t); \text{for all } y_0 \in B_j\}$ . It is easy to see that

$$\text{meas} B_j \leq C \text{meas} B_j^\varepsilon(t), \quad \text{diam} B_j^\varepsilon(t) \leq Ch.$$

Because the distance between  $X_j^\delta(t)$  and  $X_j^\varepsilon(t)$  is less than  $\|e(t)\|_\infty$ , if  $X_j^\delta(t) \in K$ , then  $B_j(t)$  lies in a disk with center in  $K$  and radius  $\delta + \|e(t)\|_\infty + Ch$ . Since  $B_j(t)$  do not overlap each other, we have

$$(3.17) \quad \text{card}\{j \in J_1; X_j^\delta(t) \in K\} \leq C \frac{\pi(\delta + \|e(t)\|_\infty + Ch)^2}{h^2}.$$

Secondly, let us estimate  $\text{card}\{j \in J_1; (X_j^\delta(t))_\delta^{(i)} \in K, X_j^\delta(t) \notin \bar{\Omega}^\delta\}$ . If  $(X_j^\delta(t))_\delta^{(i)} \in K$ , then similar to (3.15) we have

$$|(X_j^\delta(t))_\delta^{(i)} - (X_j^\varepsilon(t))^{(i)}| \leq |X_j^\varepsilon(t) - X_j^\delta(t)| + C\delta^2.$$

Let  $x_0$  be any point on  $K \cap \Omega$ . Then

$$|(X_j^\varepsilon(t))^{(i)} - x_0| \leq \delta + |X_j^\varepsilon(t) - X_j^\delta(t)| + C\delta^2.$$

If  $d$  is small enough, then the correspondence  $x \rightarrow x^{(i)}$  is one-to-one. Let  $y_0 \notin \bar{\Omega}$  and  $x_0 = (y_0)^{(i)}$ . Then

$$|X_j^\varepsilon(t) - y_0| \leq C(\delta + |X_j^\varepsilon(t) - X_j^\delta(t)| + \delta^2).$$

By the same argument we obtain an analogue of (3.17):

$$(3.18) \quad \text{card}\{j \in J_1; (X_j^\delta(t))_\delta^{(i)} \in K, X_j^\delta(t) \notin \bar{\Omega}^\delta\} \leq C \frac{(\delta + \|e(t)\|_\infty + h)^2}{h^2}.$$

Inequalities (3.16)–(3.18) imply that

$$(3.19) \quad \|e(t)\|_p \leq C\delta^2 + C \int_0^t \|e(s)\|_p ds + Ch^{\frac{2}{p}} \cdot \int_0^t \left( \frac{\delta}{h} + \frac{1}{h} \|e(s)\|_\infty + 1 \right)^{\frac{2}{p}} \delta^{-\frac{2}{p}} \|I\psi^\varepsilon(\cdot, s) - \psi^\delta(\cdot, s)\|_{1,p,\Omega^\delta} ds.$$

Using the interpolation theorem and Lemma 1, we obtain

$$(3.20) \quad \begin{aligned} & \|I\psi^\varepsilon(\cdot, s) - \psi^\delta(\cdot, s)\|_{1,p,\Omega^\delta} \\ & \leq \|\psi^\varepsilon(\cdot, s) - \psi^\delta(\cdot, s)\|_{1,p,\Omega^\delta} + \|I\psi^\varepsilon(\cdot, s) - \psi^\varepsilon(\cdot, s)\|_{1,p,\Omega^\delta} \\ & \leq C\delta^2 + C \left\{ \left( 1 + \frac{1}{\varepsilon} \|e(s)\|_\infty \right)^{\frac{2}{p}} \|e(s)\|_p + \int_0^s \|e(\tau)\|_p d\tau \right\} \\ & \quad + C\delta^2 |\psi^\varepsilon(\cdot, s)|_{3,p,\Omega^\delta}. \end{aligned}$$

By the Corollary of Theorem 1,  $|\psi^\varepsilon(\cdot, s)|_{3,p,\Omega^\delta}$  is bounded, thus substituting (3.20) into (3.19) gives (3.14).  $\square$

**THEOREM 2.** *If the assumptions of Theorem 1 hold with  $k \geq 3$  and  $d_0$  is small enough, and if  $\delta \leq C_3\varepsilon$  and there are constants  $b > 0$  and  $C_5 > 0$  such that*

$$C_5^{-1}\delta^b \leq h \leq C_5\delta,$$

then

$$(3.21) \quad \|e(t)\|_p + \|u^\varepsilon(\cdot, t) - u^\delta(\cdot, t)\|_{0,p,\Omega^\delta} \leq C\delta^2$$

for any  $p \in [1, \infty)$ .

*Proof.* It will suffice to prove the conclusion for large number  $p$ . If  $\|e(t)\|_p \leq C\delta^2$ , then the factor in (3.14),

$$\begin{aligned} 1 + \frac{1}{\delta} \|e(s)\|_\infty + \frac{h}{\delta} & \leq 1 + \frac{\|e(s)\|_p}{\delta h^{\frac{2}{p}}} + \frac{h}{\delta} \\ & \leq 1 + C\delta^{1-\frac{2b}{p}} + C_5 \\ & \leq C, \end{aligned}$$

provided  $1 - \frac{2b}{p} > 0$ . Since  $\delta \leq C_3\varepsilon$ , the factor  $(1 + \frac{1}{\varepsilon}\|e(t)\|_\infty)^{\frac{2}{q}}$  is bounded too. We have  $\|e(0)\|_p = 0$ . Thus by using Gronwall inequality and a continuous argument it is easy to prove that  $\|e(t)\|_p \leq C\delta^2$  holds for  $t \in [0, T]$ . Finally the proof of (3.21) is complete by using Lemma 1.  $\square$

**4. Full discretization.** For simplicity we assume  $f = 0$  in this section. The forward Euler scheme is applied to the ordinary differential equations (3.3). The full discretization scheme for solving  $u^{\Delta t}$ ,  $\omega^{\Delta t}$ , and  $X_j^{\Delta t}$  is the following:

$$(4.1) \quad \omega^{\Delta t}(x, n \Delta t) = \sum_{j \in J_1} \alpha_j \zeta_\varepsilon(x - X_j^{\Delta t}(n \Delta t)),$$

$$(4.2) \quad X_j^{\Delta t}((n + 1) \Delta t) = X_j^{\Delta t}(n \Delta t) + \Delta t g^{\Delta t}(X_j^{\Delta t}(n \Delta t), n \Delta t),$$

$$(4.3) \quad X_j^{\Delta t}(0) = X_j,$$

$\psi^{\Delta t}(n \Delta t) \in V^\delta$ , and

$$(4.4) \quad \int_{\Omega^\delta} \nabla \psi^{\Delta t} \cdot \nabla v dx = \int_{\Omega^\delta} \omega^{\Delta t} v dx \quad \forall v \in V^\delta,$$

$$(4.5) \quad u^{\Delta t}(n \Delta t) = \nabla \wedge \psi^{\Delta t}(n \Delta t),$$

where  $\Delta t$  is the length of time step, and

$$g^{\Delta t}(x, t) = \sum_{i=1}^M a_i u^{\Delta t}(x_\delta^{(i)}, t).$$

Now we estimate the error  $X_j - X_j^{\Delta t}$  and  $u^\varepsilon - u^{\Delta t}$ . Let

$$\|e^n\|_p = \left( h^2 \sum_{j \in J_3} |X_j^{\Delta t}(n \Delta t) - X_j(n \Delta t)|^p \right)^{\frac{1}{p}},$$

$$\|e^n\|_\infty = \max_{j \in J_3} |X_j^{\Delta t}(n \Delta t) - X_j(n \Delta t)|,$$

where

$$J_3 = \left\{ j; \quad X_j \in \Omega_d \cap \Omega_{2C_0(\varepsilon + \delta + \Delta t \frac{1}{2})} \right\}.$$

LEMMA 3. *There exists a constant  $C_6$  such that*

$$(4.6) \quad |X((n + 1) \Delta t) - X(n \Delta t) - \Delta t \frac{dX}{dt}(n \Delta t)| \leq C_6 \Delta t^2.$$

*Proof.* Since  $u$  is sufficiently smooth, (4.6) obviously follows.  $\square$

LEMMA 4. *Under the assumptions of Theorem 1,*

$$(4.7) \quad \left( h^2 \sum_{j \in J_3} |u^\varepsilon((X_j^{\Delta t}(n \Delta t))_\delta^{(i)}, n \Delta t) - u((X_j(n \Delta t))_\delta^{(i)}, n \Delta t)|^p \right)^{\frac{1}{p}}$$

$$\leq C(\varepsilon^k + \delta^2 + \|e^n\|_p),$$

where  $p \geq 2$ .

*Proof.* We have

$$\begin{aligned} & |u^\varepsilon((X_j^{\Delta t}(n \Delta t))_\delta^{(i)}, n \Delta t) - u((X_j(n \Delta t))_\delta^{(i)}, n \Delta t)| \\ & \leq |u^\varepsilon((X_j^{\Delta t}(n \Delta t))_\delta^{(i)}, n \Delta t) - u^\varepsilon((X_j(n \Delta t))_\delta^{(i)}, n \Delta t)| \\ & \quad + |u^\varepsilon((X_j(n \Delta t))_\delta^{(i)}, n \Delta t) - u^\varepsilon((X_j(n \Delta t))^{(i)}, n \Delta t)| \\ & \quad + |u^\varepsilon((X_j(n \Delta t))^{(i)}, n \Delta t) - u((X_j(n \Delta t))^{(i)}, n \Delta t)| \\ & \quad + |u((X_j(n \Delta t))^{(i)}, n \Delta t) - u((X_j(n \Delta t))_\delta^{(i)}, n \Delta t)|. \end{aligned}$$

By (2.17) and the Appendix, we get

$$\begin{aligned} & |u^\varepsilon((X_j^{\Delta t}(n \Delta t))_\delta^{(i)}, n \Delta t) - u^\varepsilon((X_j(n \Delta t))_\delta^{(i)}, n \Delta t)| \\ & \leq C_1 |(X_j^{\Delta t}(n \Delta t))_\delta^{(i)} - (X_j(n \Delta t))_\delta^{(i)}| \\ & \leq C_1 (|(X_j^{\Delta t}(n \Delta t))_\delta^{(i)} - (X_j^{\Delta t}(n \Delta t))^{(i)}| \\ & \quad + |(X_j^{\Delta t}(n \Delta t))^{(i)} - (X_j(n \Delta t))^{(i)}| + |(X_j(n \Delta t))^{(i)} - (X_j(n \Delta t))_\delta^{(i)}|) \\ & \leq C\delta^2 + C|X_j^{\Delta t}(n \Delta t) - X_j(n \Delta t)|, \end{aligned}$$

$$\begin{aligned} & |u^\varepsilon((X_j(n \Delta t))_\delta^{(i)}, n \Delta t) - u^\varepsilon((X_j(n \Delta t))^{(i)}, n \Delta t)| \\ & \leq C|(X_j(n \Delta t))_\delta^{(i)} - (X_j(n \Delta t))^{(i)}| \\ & \leq C\delta^2, \end{aligned}$$

and

$$\begin{aligned} & |u((X_j(n \Delta t))^{(i)}, n \Delta t) - u((X_j(n \Delta t))_\delta^{(i)}, n \Delta t)| \\ & \leq C|(X_j(n \Delta t))^{(i)} - (X_j(n \Delta t))_\delta^{(i)}| \\ & \leq C\delta^2. \end{aligned}$$

Then, we have

$$\begin{aligned} & \left( h^2 \sum_{j \in J_3} |u^\varepsilon((X_j^{\Delta t}(n \Delta t))_\delta^{(i)}, n \Delta t) - u((X_j(n \Delta t))_\delta^{(i)}, n \Delta t)|^p \right)^{\frac{1}{p}} \\ & \leq C\delta^2 + \left( h^2 \sum_{j \in J_3} |u^\varepsilon((X_j(n \Delta t))^{(i)}, n \Delta t) \right. \\ & \quad \left. - u((X_j(n \Delta t))^{(i)}, n \Delta t)|^p \right)^{\frac{1}{p}} + C\|e^n\|_p. \end{aligned}$$

By (2.25), (2.26) of [10],

$$\begin{aligned} & \left( h^2 \sum_{j \in J_3} |u^\varepsilon((X_j(n \Delta t))^{(i)}, n \Delta t) - u((X_j(n \Delta t))^{(i)}, n \Delta t)|^p \right)^{\frac{1}{p}} \\ & \leq C(\|u - u^\varepsilon\|_{0,p,\Omega} + h|u - u^\varepsilon|_{1,p,\Omega}) \\ & \leq C\varepsilon^k. \end{aligned}$$

Thus (4.7) is obtained.  $\square$

We will prove by induction that

$$(4.8) \quad \|e^l\|_p \leq C_7(\varepsilon^k + \delta^2 + \Delta t)$$

holds for a suitable  $C_7$  and all  $l$ . Since  $\|e^0\|_p = 0$ , we assume that (4.8) is valid for  $0 \leq l \leq n$ .

From now on we assume that all the hypotheses of Theorem 2 are satisfied. Since only large  $p$  is needed to be taken into consideration, we may assume the constant  $a$  in (2.16) satisfies  $a < \frac{p}{2}$ . Also we assume

$$(4.9) \quad \Delta t \leq C_8 \delta^2, \quad \varepsilon^{k-1} \leq C_8 \delta.$$

Again by (2.31) of [10] and (2.16) we have

$$(4.10) \quad \|e^l\|_\infty \leq \frac{1}{h^{\frac{2}{p}}} \|e^l\|_p \leq \frac{\tilde{C}_p^{\frac{2}{p}}}{\varepsilon^{\frac{2a}{p}}} \|e^l\|_p.$$

Then (4.8) and  $\delta \leq C_3 \varepsilon$  yield

$$\|e^l\|_\infty \leq C(\varepsilon^{k-1} + \delta + \Delta t^{\frac{1}{2}}) \varepsilon^{1 - \frac{2a}{p}}.$$

Noting  $k \geq 3$ , we get

$$(4.11) \quad \|e^l\|_\infty \leq \varepsilon + \delta + \Delta t^{\frac{1}{2}}$$

for sufficiently small  $\varepsilon$ .

LEMMA 5. *If the hypotheses of Theorem 2 and (4.9) hold, and (4.8) holds for  $l = n$ , and if  $\varepsilon$  is sufficiently small, then*

$$\begin{aligned} & \|\psi^{\Delta t}(\cdot, n \Delta t) - \psi^\varepsilon(\cdot, n \Delta t)\|_{1,p,\Omega^\delta} \leq C_9 \delta^2 \\ & + C_9 \left(1 + \frac{1}{\varepsilon} \|e^n\|_\infty\right)^{\frac{2}{q}} \left(h^2 \sum_{j \in J_3} |X_j^{\Delta t}(n \Delta t) - X_j^\varepsilon(n \Delta t)|^p\right)^{\frac{1}{p}}, \end{aligned}$$

where  $p \geq 2, \frac{1}{p} + \frac{1}{q} = 1$ , and the constant  $C_9$  is independent of  $C_7$ .

*Proof.* The proof is almost the same as that of Lemma 1. Now

$$\psi^\varepsilon - \psi^\delta = \phi_1 + \phi_2,$$

$$\begin{aligned} \phi_1 &= \Delta^{-1} \left( \sum_{j \in J_1} \alpha_j (\zeta_\varepsilon(\cdot - X_j^\varepsilon(n \Delta t)) - \zeta_\varepsilon(\cdot - X_j^{\Delta t}(n \Delta t))) \right), \\ \phi_2 &= (\Delta^{-1} - \Delta_\delta^{-1}) \sum_{j \in J_1} \alpha_j \zeta_\varepsilon(\cdot - X_j^{\Delta t}(n \Delta t)). \end{aligned}$$

Inequality (4.11) implies

$$\phi_1 = \Delta^{-1} \left( \sum_{j \in J_3} \alpha_j (\zeta_\varepsilon(\cdot - X_j^\varepsilon(n \Delta t)) - \zeta_\varepsilon(\cdot - X_j^{\Delta t}(n \Delta t))) \right).$$

Using the same argument we can estimate  $\phi_1$  and  $\phi_2$ .  $\square$

LEMMA 6. *If the hypotheses of Theorem 2 and (4.9) hold, and if  $\varepsilon$  is sufficiently small, then*

$$(4.12) \quad \left( h^2 \sum_{j \in J_3} |g^{\varepsilon\delta}(X_j^{\Delta t}(n \Delta t), n \Delta t) - g^{\Delta t}(X_j^{\Delta t}(n \Delta t), n \Delta t)|^p \right)^{\frac{1}{p}} \leq C \left( 1 + \frac{1}{\delta} \|e^n\|_\infty \right) (\delta^2 + \|\psi^\varepsilon(\cdot, n \Delta t) - \psi^{\Delta t}(\cdot, n \Delta t)\|_{1,p,\Omega^\varepsilon}).$$

*Proof.* Like the proof of Lemma 2, we can get

$$(4.13) \quad \|I\psi^\varepsilon(\cdot, t) - \psi^\varepsilon(\cdot, t)\|_{1,\infty,\Omega^\varepsilon} \leq C\delta^2 |\psi^\varepsilon(\cdot, t)|_{3,\infty,\Omega} \leq C\delta^2,$$

$$(4.14) \quad \|I\psi^\varepsilon(\cdot, t) - \psi^{\Delta t}(\cdot, t)\|_{1,\infty,K_j^{(\varepsilon)}} \leq C\delta^{-\frac{2}{p}} \|I\psi^\varepsilon(\cdot, t) - \psi^{\Delta t}(\cdot, t)\|_{1,p,K_j^{(\varepsilon)}}.$$

Then

$$\begin{aligned} & |g^{\Delta t}(X_j^{\Delta t}(n \Delta t), n \Delta t) - g^{\varepsilon\delta}(X_j^{\Delta t}(n \Delta t), n \Delta t)| \\ & \leq \sum_{i=1}^M |a_i| \cdot \|u^\varepsilon - u^{\Delta t}\|_{0,\infty,K_j^{(\varepsilon)}} \\ & = \sum_{i=1}^M |a_i| \cdot \|\psi^\varepsilon - \psi^{\Delta t}\|_{1,\infty,K_j^{(\varepsilon)}} \\ & \leq C\delta^2 + C\delta^{-\frac{2}{p}} \sum_{i=1}^M |a_i| \cdot \|I\psi^\varepsilon(\cdot, n \Delta t) - \psi^{\Delta t}(\cdot, n \Delta t)\|_{1,p,K_j^{(\varepsilon)}}. \end{aligned}$$

For any  $K \in \mathcal{T}_\delta$ , we also have

$$\text{card}\{j \in J_1, X_j^{\Delta t}(n \Delta t) \in K\} \leq C \frac{(\delta + \|e^n\|_\infty + h)^2}{h^2},$$

and therefore

$$\begin{aligned} & \left( h^2 \sum_{j \in J_3} |g^{\Delta t}(X_j^{\Delta t}(n \Delta t), n \Delta t) - g^{\varepsilon\delta}(X_j^{\Delta t}(n \Delta t), n \Delta t)|^p \right)^{\frac{1}{p}} \\ & \leq C\delta^2 + C \left( 1 + \frac{1}{\delta} \|e^n\|_\infty \right)^{\frac{2}{p}} \|I\psi^\varepsilon(\cdot, n \Delta t) - \psi^{\Delta t}(\cdot, n \Delta t)\|_{1,p,\Omega^\varepsilon}. \end{aligned}$$

In view of (4.13) we obtain (4.12).  $\square$

LEMMA 7. *Under the hypotheses of Lemma 5,*

$$(4.15) \quad \|e^{n+1}\|_p \leq \|e^n\|_p + C_{10} \Delta t \left( 1 + \frac{1}{\varepsilon} \|e^n\|_\infty \right)^{\frac{2}{q}} \left( 1 + \frac{1}{\delta} \|e^n\|_\infty \right)^{\frac{1}{p}} \cdot (\varepsilon^k + \delta^2 + \Delta t + \|e^n\|_p),$$

where the constant  $C_{10}$  is independent of  $C_7$ .

*Proof.* By Lemma 3 and (4.2) we obtain

$$\begin{aligned}
 \|e^{n+1}\|_p &= \left( h^2 \sum_{j \in J_3} |X_j^{\Delta t}((n+1)\Delta t) - X_j((n+1)\Delta t)|^p \right)^{\frac{1}{p}} \\
 (4.16) \quad &\leq \|e^n\|_p + C \Delta t^2 + \Delta t \left( h^2 \sum_{j \in J_3} |g^{\Delta t}(X_j^{\Delta t}(n\Delta t), n\Delta t) \right. \\
 &\quad \left. - u(X_j(n\Delta t), n\Delta t)|^p \right)^{\frac{1}{p}}.
 \end{aligned}$$

We have

$$\left( h^2 \sum_{j \in J_3} |g^{\Delta t}(X_j^{\Delta t}(n\Delta t), n\Delta t) - u(X_j(n\Delta t), n\Delta t)|^p \right)^{\frac{1}{p}} \leq I_1 + I_2,$$

where

$$\begin{aligned}
 I_1 &= \left( h^2 \sum_{j \in J_3} |g^{\Delta t}(X_j^{\Delta t}(n\Delta t), n\Delta t) - g^{\varepsilon\delta}(X_j^{\Delta t}(n\Delta t), n\Delta t)|^p \right)^{\frac{1}{p}}, \\
 I_2 &= \left( h^2 \sum_{j \in J_3} |g^{\varepsilon\delta}(X_j^{\Delta t}(n\Delta t), n\Delta t) - u(X_j(n\Delta t), n\Delta t)|^p \right)^{\frac{1}{p}},
 \end{aligned}$$

$$\begin{aligned}
 I_2 &= \left( h^2 \sum_{j \in J_3} \left| \sum_{i=1}^M a_i (u^\varepsilon((X_j^{\Delta t}(n\Delta t))_\delta^{(i)}, n\Delta t) - u(X_j(n\Delta t), n\Delta t)) \right|^p \right)^{\frac{1}{p}} \\
 &\leq \left( h^2 \sum_{j \in J_3} \left| \sum_{i=1}^M a_i (u^\varepsilon((X_j^{\Delta t}(n\Delta t))_\delta^{(i)}, n\Delta t) - u((X_j(n\Delta t))_\delta^{(i)}, n\Delta t)) \right|^p \right)^{\frac{1}{p}} \\
 &\quad + \left( h^2 \sum_{j \in J_3} \left| \sum_{i=1}^M a_i (u((X_j(n\Delta t))_\delta^{(i)}, n\Delta t) - u((X_j(n\Delta t))^{(i)}, n\Delta t)) \right|^p \right)^{\frac{1}{p}} \\
 &\quad + \left( h^2 \sum_{j \in J_3} \left| \sum_{i=1}^M a_i (u((X_j(n\Delta t))^{(i)}, n\Delta t) - u((X_j(n\Delta t)), n\Delta t)) \right|^p \right)^{\frac{1}{p}} \\
 &\leq \sum_{i=1}^M |a_i| \left( h^2 \sum_{j \in J_3} |u^\varepsilon((X_j^{\Delta t}(n\Delta t))_\delta^{(i)}, n\Delta t) - u((X_j(n\Delta t))_\delta^{(i)}, n\Delta t)|^p \right)^{\frac{1}{p}} \\
 &\quad + \sum_{i=1}^M |a_i| \left( h^2 \sum_{j \in J_3} |u((X_j(n\Delta t))_\delta^{(i)}, n\Delta t) - u((X_j(n\Delta t))^{(i)}, n\Delta t)|^p \right)^{\frac{1}{p}} \\
 &\quad + \left( h^2 \sum_{j \in J_3} \left| \sum_{i=1}^M a_i (u((X_j(n\Delta t))^{(i)}, n\Delta t) - u(X_j(n\Delta t), n\Delta t)) \right|^p \right)^{\frac{1}{p}}
 \end{aligned}$$

Using Lemma 4, the fact that  $u$  is smooth enough, and Taylor’s formula, we get

$$I_2 \leq C(\varepsilon^k + \delta^2 + \|e^n\|_p).$$

The term in (4.16) with respect to  $I_1$  can be estimated by using Lemma 6. Then we get

$$(4.17) \quad \begin{aligned} \|e^{n+1}\|_p &\leq \|e^n\|_p + C \Delta t^2 + C \Delta t \left(1 + \frac{1}{\delta} \|e^n\|_\infty\right)^{\frac{2}{p}} \\ &\quad \cdot (\delta^2 + \|\psi^\varepsilon(\cdot, n \Delta t) - \psi^{\Delta t}(\cdot, n \Delta t)\|_{1,p,\Omega^\delta}) \\ &\quad + C \Delta t(\varepsilon^k + \delta^2 + \|e^n\|_p). \end{aligned}$$

By Theorem 1

$$\left(h^2 \sum_{j \in J_3} |X_j(t) - X_j^\varepsilon(t)|^p\right)^{\frac{1}{p}} \leq C_2(2(\varepsilon + \delta + \Delta t^{\frac{1}{2}}))^k \leq C\varepsilon^k.$$

Hence Lemma 5 implies

$$\|\psi^{\Delta t}(\cdot, n \Delta t) - \psi^\varepsilon(\cdot, n \Delta t)\|_{1,p,\Omega^\delta} \leq C\delta^2 + C \left(1 + \frac{1}{\varepsilon} \|e^n\|_\infty\right)^{\frac{2}{q}} (\varepsilon^k + \|e^n\|_p).$$

Substituting this into (4.17) gives (4.15).  $\square$

**THEOREM 3.** *If the hypotheses of Theorem 2 and (4.9) hold, then*

$$(4.18) \quad \|e^n\|_p + \|u^\varepsilon(\cdot, n \Delta t) - u^{\Delta t}(\cdot, n \Delta t)\|_{0,p,\Omega^\delta} \leq C(\varepsilon^k + \delta^2 + \Delta t),$$

where  $n \Delta t \leq T$ .

*Proof.* It will suffice to prove (4.18) for large  $p$  and small  $\varepsilon$ . We have assumed that (4.8) is valid for  $0 \leq l \leq n$ . From (4.10) and the relations among  $\varepsilon, \delta$ , and  $\Delta t$ , we get

$$\frac{1}{\delta} \|e^n\|_\infty \leq C_7 \tilde{C}^{\frac{2}{p}} (\varepsilon^k + \delta^2 + \Delta t) \varepsilon^{-\frac{2a}{p}} \delta^{-1} \leq C_{11},$$

for  $\varepsilon \leq \varepsilon_0$ , where  $\varepsilon_0$  depends on  $C_7$  but  $C_{11}$  is independent of  $C_7$ . Analogously we can estimate  $\frac{1}{\varepsilon} \|e^n\|_\infty$ . Then (4.15) becomes

$$\|e^{n+1}\|_p \leq \|e^n\|_p + C_{12} \Delta t(\varepsilon^k + \delta^2 + \Delta t + \|e^n\|_p),$$

for  $\varepsilon \leq \varepsilon_0$ , where  $C_{12}$  is independent of  $C_7$ .

We set  $C_7 = C_{12} T e^{C_{12} T}$ . Then we determine  $\varepsilon_0$  according to  $C_7$ . It is easy to verify that

$$\|e^n\|_p \leq C_{12} e^{C_{12} n \Delta t} (\varepsilon^k + \delta^2 + \Delta t) n \Delta t$$

for  $\varepsilon \leq \varepsilon_0$  and all  $n, n \Delta t \leq T$ . Thus the estimate for  $\|e^n\|_p$  is obtained, and the estimate for  $u^\varepsilon - u^{\Delta t}$  follows from Lemma 5.  $\square$

*Remark.* The extension of convergence proof to higher order schemes for time stepping is straightforward. A numerical example was given to show the accuracy of this method [11].

**Appendix.** Let domains  $\Omega$  and  $\Omega^\delta$  be the above. We prove

$$(A.1) \quad \text{meas}(\Omega^\delta \setminus \Omega) \leq C\delta^3$$

and

$$(A.2) \quad |x^{(i)} - x_\delta^{(i)}| \leq C\delta^2 \quad \forall x \in \Omega_d,$$

provided  $d$  is small enough.

Introducing local coordinates,  $\partial\Omega^\delta$  is the quadratic interpolation of  $\partial\Omega$ . Thus we have (Chapter 2 of [9], for example)

$$\sup_{x \in \partial\Omega^\delta} \inf_{y \in \partial\Omega} |x - y| \leq C\delta^3,$$

from which (A.1) follows.

Let us consider (A.2). If  $x \in \bar{\Omega}$ , then (A.2) is trivial, so we assume  $x \notin \bar{\Omega}$ . From Fig. 3.1 it is clear that

$$|a_{12,K} - \tilde{a}_{12,K}| \leq C\delta^2.$$

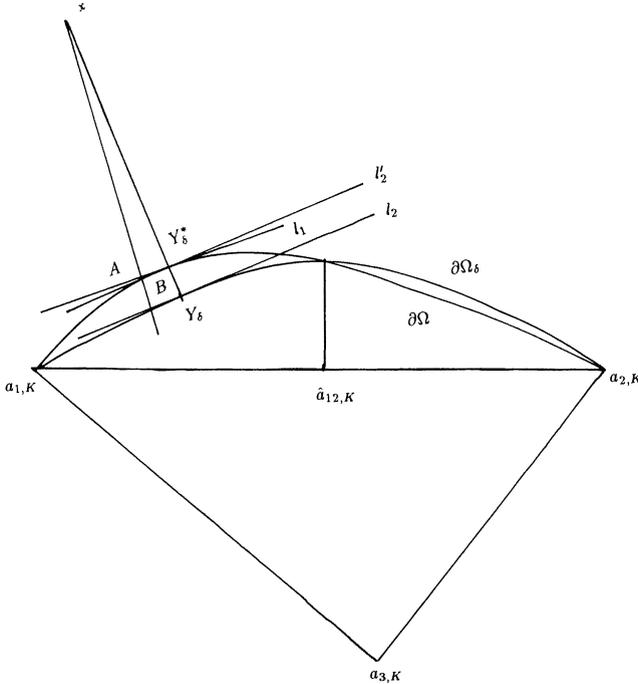


FIG. A.1

Let  $Y_\delta^*$  be the intersecting point of line  $\overline{xY_\delta^*}$  and  $\partial\Omega$ ,  $l_1$  be the tangent line of  $\partial\Omega$  through  $Y_\delta^*$ ,  $l_2$  be the tangent line of  $\partial\Omega^\delta$  through  $Y_\delta$ , and  $l_2'$  be the parallel line of  $l_2$  through  $Y_\delta^*$  (Fig. A.1). The angle between  $l_1$  and  $l_2'$  is less than  $C\delta^2$ . We draw the vertical line of  $l_1$  through point  $x$ . Let  $A$  be the foot of perpendicular and  $B$  be

the intersecting point of  $\overline{x\bar{A}}$  with  $l'_2$ . Then  $Y$  lies in the triangle  $Y_\delta^*AB$ . The angle  $\angle Y_\delta^*xB$  is equal to the angle between  $l_1$  and  $l'_2$ . Now  $x \in \Omega_d$ , hence  $|xY_\delta^*| \leq Cd$ . Then

$$|Y_\delta^*Y| \leq |Y_\delta^*B| \leq Cd \cdot C\delta^2 \leq C\delta^2.$$

Consequently,

$$|Y_\delta Y| \leq |Y_\delta Y_\delta^*| + |Y_\delta^*Y| \leq C\delta^3 + C\delta^2 \leq C\delta^2.$$

By definition,

$$|x_\delta^{(i)} - x^{(i)}| = (i+1)|Y_\delta Y| \leq C\delta^2,$$

which proves (A.2) if  $Y_\delta$  is not a node. Conversely, if  $Y_\delta$  is just a node, say  $a_{1,K}$ , then we consider the two triangles containing  $a_{1,K}$ . The argument is similar.

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